

Spin-orbit splitting of valence and conduction bands in HgTe quantum wells near the Dirac point

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Energy spectra both of the conduction and valence bands of the HgTe quantum wells with a width close to the Dirac point were studied experimentally. Simultaneous analysis of the Shubnikov-de Haas oscillations and Hall effect over a wide range of electron and hole densities gives surprising result: the top of the valence band is strongly split by spin-orbit interaction while the splitting of the conduction band is absent, within experimental accuracy. Astonishingly, but such a ratio of the splitting values is observed as for structures with normal spectrum so for structures with inverted one. These results do not consistent with the results of kP calculations, in which the smooth electric field across the quantum well is only reckoned in. It is shown that taking into account the asymmetry of the quantum well interfaces within a tight-binding method gives reasonable agreement with the experimental data.

I. INTRODUCTION

Heightened interest to two-dimensional (2D) structures with quantum well of gapless semiconductors is caused by the fact that different types of the carrier energy spectrum could be realized in these structures depending on the width (d) of quantum well. The most detailed, both theoretically and experimentally, heterostructures with $\text{Hg}_{1-x}\text{Cd}_x\text{Te}/\text{HgTe}/\text{Hg}_{1-x}\text{Cd}_x\text{Te}$ quantum well were studied. The calculations of the energy spectrum in framework of kP method for symmetrical quantum well show that there is the critical width of the HgTe quantum well, $d = d_c \simeq 6.3$ nm, when the linear in quasimomentum Dirac-like energy spectrum should be realized at small quasimomentum (k) [1, 2]. At $d < d_c$ the energy spectrum is normal. The valence band is formed from heavy-hole states while the conduction band is formed from electron states and states of light hole. At $d > d_c$ these states swap over, and such a spectrum is called an inverted spectrum. Another approach, the calculation in the framework of tight-binding model, which takes into account the bulk inversion asymmetry of the zinc blende lattice, gives close result regarding the dependence $E(k)$ except the fact that the relatively large anti-crossing at $k = 0$ and large spin-orbit (SO) splitting of the conduction and valence band appears at $d = d_c$.

Knowledge of the energy spectrum is necessary for an understanding of all properties of 2D systems: optical, transport and others. However it has been experimentally studied rather superficially to date. The energy spectrum of the conduction band was studied more detailed [3–7]. There was shown that the SO splitting of the conduction band in the technologically symmetric quantum wells was not reveal itself in the most cases, the effective mass of the electrons increased with their density. These data were found in satisfactory agreement

with results of theoretical calculations performed in the framework of the kP method.

The energy spectrum of the valence band is investigated still less and, what is important, the experimental data are inconsistent with theoretical results in many cases. So in the structures with $d > 10$ nm that corresponds to the inverted energy spectrum, the hole effective mass (m_h) within wide hole density range $p = (1-4) \times 10^{11} \text{ cm}^{-2}$ occurs substantially less than that calculated within kP method: $m_h \simeq (0.15 - 0.3)m_0$ [8–10] instead of $(0.5 - 0.6)m_0$ [11]. These calculations predict also that the conduction band should overlap with the valence band which top is located at $k \neq 0$ when $d \gtrsim (12 - 15)$ nm. Although this prediction is in qualitative agreement with experimental data [10, 12], the quantitative difference between theory and experiment is drastic. Experimentally, the top of the valence band is located at $k \simeq 0.5 \times 10^6 \text{ cm}^{-1}$, while the theoretical prediction gives the value of about $k \simeq 2.5 \times 10^6 \text{ cm}^{-1}$.

The valence band spectrum in the heterostructures with normal band ordering was studied in several papers only. The experimental data published in Refs. [13] and [14] for structures with $d = (5 - 6)$ nm are very similar however the interpretation differs significantly. Analyzing the Hall density and the Fourier spectra of the Shubnikov-de Haas (SdH) oscillations, the authors of Ref. [13] are able to describe the data taking into account the secondary maxima of the dispersion law located at $k \neq 0$. But it has been done only for one hole density. Studying analogous heterostructures, the authors of Ref. [14] show that such interpretation does not describe the experimental data within a wide hole density range. They show that all the results are well described under assumption that the top of the valence band is very strongly split by SO interaction. They supposed that the reason for such a splitting is a strong electric field of p - n junction of technological origin in which the quantum is

embedded. The authors were unable to investigate the splitting of the conduction band. These data would make the interpretation more reliable.

In the present paper, we report the results of an experimental study both of hole and electron transport in HgTe quantum well of different width near the critical point d_c with a normal and inverted energy spectra. The measurements were performed over a wide range of carrier density. It has been experimentally found that the valence band is strongly split by SO interaction, while the conduction band remains unsplit independently of energy band ordering. We believe that the natural interface inversion asymmetry of zinc-blende heterostructures is responsible for these peculiarities. Quantitatively, the experimental data are well described in the framework of atomistic calculation [15].

II. EXPERIMENTAL

Our samples with HgTe quantum wells were realized on the basis of HgTe/Hg_{1-x}Cd_xTe ($x = 0.55 - 0.65$) heterostructures grown by molecular beam epitaxy on GaAs substrate with the (013) surface orientation [16]. The samples were mesa etched into standard Hall bars of 0.5 mm width and the distance between the potential probes of 0.5 mm. To change and control the carrier density in the quantum well, the field-effect transistors were fabricated with parylene as an insulator and aluminium as a gate electrode. For each heterostructure, several samples were fabricated and studied. The measurements were performed at temperatures of 1.3 – 20 K.

III. CHARACTERIZATION OF SAMPLES REGARDING THE TYPE OF SPECTRUM

It is very important for interpretation of the experimental results to know whether the spectrum of structure under study is normal or inverted. The values of quantum well width presented in the Table I are technological, therefore it is very desirable to have independent data on the spectrum type.

The most reliable method to determine the type of the spectra can be based on the peculiarity of the spectrum quantization in the external magnetic field. Theoretical calculations [5, 11, 17] and experimental investigations [10, 17] show that there are two anomalous Landau levels which behavior is radically different for the normal and inverted spectra [Figs. 1(c) and 1(f)].

As seen from Fig. 1(f) the Landau levels -2 and 0 in the structure with inverted spectrum start at $B = 0$ from the bottom of the conduction band and top of the valence, respectively, and moving towards each other cross in the magnetic field $B = B_c$. This magnetic field increases with d increase, achieves maximal value of about 9–10 T at $d \simeq (9-10)$ nm and then decreases with a further d increase. In the structures with normal spectrum,

Table I. The parameters of heterostructures under study

#	d (nm)	p_s (cm ⁻²)	B_c (T)	Q/e (cm ⁻²) ^a	p_1/p_2
1122	5.6	1.3×10^{11}	–	2.5×10^{11}	1.8 – 2.2
H724	5.8	1.5×10^{11}	–	1.0×10^{11}	2.0 – 2.2
1123	6.0	6.0×10^{10}	–	1.5×10^{11}	2.0 – 2.3
1121	6.3	8.0×10^{10}	0.6 ± 0.2	7.6×10^{10}	2.2 – 2.5
1023	6.5	3.6×10^{10}	1.1 ± 0.2	1.0×10^{10} -7.0×10^{10b}	2.3 – 2.7
1022	6.7	1.4×10^{10}	1.2 ± 0.2	1.5×10^{10} -1.5×10^{10b}	2.3 – 2.8
1124	7.1	3.4×10^{10}	~ 0	2.2×10^{11}	2.2 – 2.8

^a At $V_g = 0$.

^b After illumination.

the energy positions of the anomalous Landau levels at $B \rightarrow 0$ is opposite: the level -2 starts from the valence band top while the level 0 goes from the conduction band bottom. Because they move in the same directions as for $d > d_c$, the crossing does not occur in this case [Fig. 1(c)]. Thus, if the crossing of Landau levels is detected experimentally, the real width of the HgTe quantum well is greater than critical value d_c and the sample under study is in inverted regime.

When B_c is large enough so that the anomalous Landau levels have a large density of states and are well separated, the cross manifests itself as the nonmonotonic peculiarity in the ρ_{xx} vs B and ρ_{xy} vs B dependences [10]. When the quantum well width is close to the critical value d_c , the cross of the Landau levels occurs at so low magnetic field that it may not reveal itself in magnetotransport measurements. It is possible in this case to observe the cross by studying the behavior of the quantum capacitance (C_q) in the magnetic field at the gate voltages close to the charge neutrality point (CNP). Unlike the resistance components which depend not only on the density of states (ν) but on the disorder strength also, the quantum capacitance depends on the density of states only, $C_q = e^2 dn/d\mu = e^2 \nu(\mu)$, where μ stands for the chemical potential. So, the value of C_q at $V_g \simeq V_g^{\text{CNP}}$ should increase with growing B , achieve the maximal value at $B = B_c$ and, then, decrease with the further increase of magnetic field even when the Landau levels are rather broadened. When the Fermi level is located in the bands at $V_g \neq V_g^{\text{CNP}}$, the maximum in the dependence $C_q(B)$ should be also observed but in the higher magnetic fields as compared with B_c .

In the structures with normal energy spectrum, the anomalous Landau levels are moving apart and the density of states in the gap (which is nonzero due to smearing) decreases with growing magnetic field. Thus, the capacitance should decrease with B near CNP.

The measurements of the quantum capacitance were taken in all the structures investigated. In Fig. 2 we have presented, as an example, the dependence $C_q(V_g)$

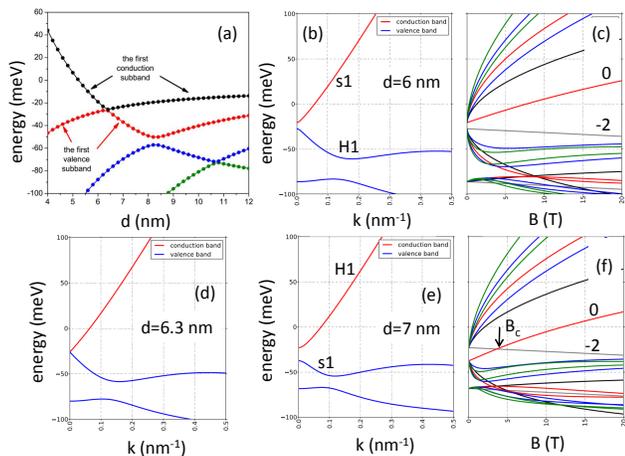


Figure 1. (Color online) The spatial quantization subband energies at $k = 0$ plotted against the HgTe quantum well width. The band structure (b), (d), (e) and Landau levels (c), (f) for quantum well of different widths [11].

measured at $B = 0$ on the structure 1023. The non-monotonic volt-capacitance characteristic is consequence of the nonmonotonic energy dependence of the density of states. The minimum in the vicinity of CNP is observed when the Fermi level goes through the energy gap where the density of states is much smaller than that in the valence and conduction bands [18]. The upper inset in Fig. 2 shows the magnetic field dependences of C_q at different V_g near CNP. The maximum in these curves at $B = B_{max}$, which shifts to the higher B values when V_g deviates from CNP, is clearly evident. As seen from the lower inset in Fig. 2 the B_{max} vs V_g plot has a minimum and just the minimal value of $B_{max} \simeq 1.2$ T corresponds to the magnetic field B_c in which the anomalous Landau levels 0 and -2 cross each other. The analogous maximum in the dependence $C_q(B)$ at V_g close to V_g^{CNP} was also observed in the structures 1022 and 1121. The B_c values obtained in such a manner for these three structures are listed in the Table I.

As for the remaining four heterostructures, the quantum capacitance for the structure 1124 in the vicinity of CNP is practically independent on B , while for the structures 1122, H724, and 1123 C_q decreases when B increases.

Thus the results of capacitance measurements in the presence of magnetic field show unambiguously that the structures 1022, 1023 and 1121 has inverted spectrum, the structure 1124 is very close to the critical point $d = d_c$, while the other structures have normal spectra.

IV. THE VALENCE BAND SPECTRUM

Let us now inspect the gate voltage dependences of the Hall carrier density obtained at low magnetic field [where $R_H(B)$ is constant] as follows: $n_H =$

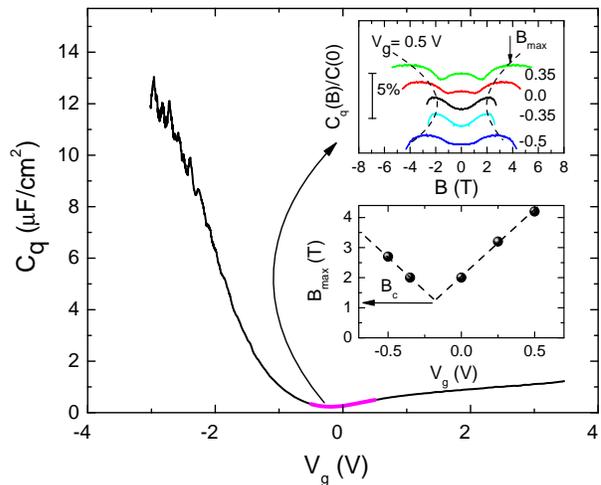


Figure 2. (Color online) The gate voltage dependence of the quantum capacitance C_q for the structure 1023 at $T = 4.2$ K and $B = 0$. The upper inset is the magnetic field dependences of C_q at the gate voltages close to CNP. The lower inset shows the B_{max} values plotted against the gate voltage.

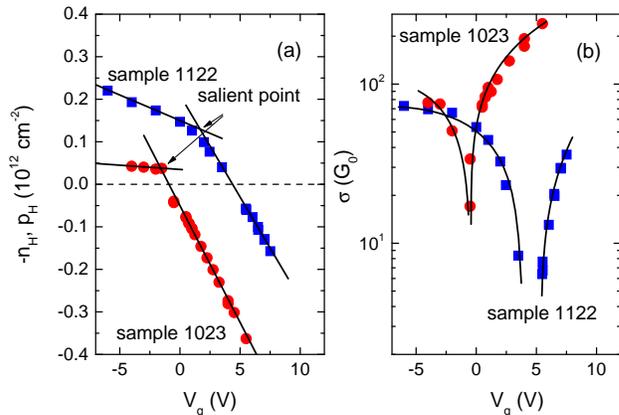


Figure 3. (Color online) The gate voltage dependences of (a) the Hall carrier densities n_H and p_H and (b) the conductivity at $B = 0$ for the structures 1122 and 1023 with $d < d_c$ and $d > d_c$, respectively. $T = 4.2$ K.

$-1/eR_H(0.1 \text{ T})$ when $R_H < 0$ for electron density and $p_H = 1/eR_H(0.1 \text{ T})$ when $R_H > 0$ for hole density. These dependences are plotted in Fig. 3(a) for structures 1122 and 1023 with normal and inverted spectrum, respectively. The corresponding V_g dependences of the conductivity are shown in Fig. 3(b). These structures are chosen as typical ones, the results for other structures are analogous. In what follows we will demonstrate all the results namely for these structures excepting the cases when the simultaneous analysis of results for all structures is useful. One can see that the electron density depends linearly on V_g over the whole gate voltage range. Note that the slope dn_H/dV_g coincides with the value C/eS (where C is the capacitance between the 2D gas and gate electrode and S is the gate area), to within

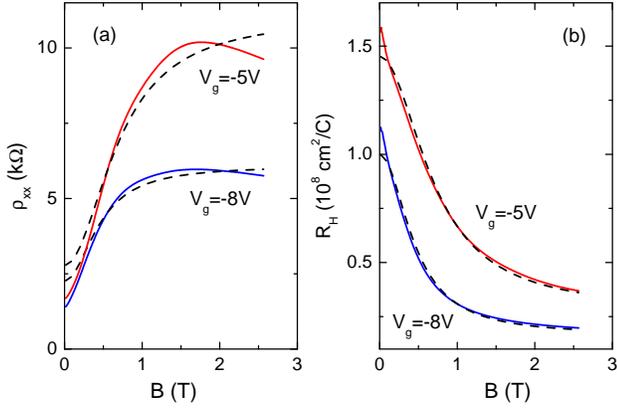


Figure 4. (Color online) The magnetic field dependences of ρ_{xx} (a) and R_H (b) for two gate voltages for the structure 1023, $T = 18$ K. The solid curves are measured experimentally, the dashed ones are the results of the best fit with the parameters shown in Fig. 5.

experimental error. The hole density changes with the same slope, $dp_H/dV_g = -dn_H/dV_g$, within a restricted V_g range only. At the gate voltage $V_g = V_g^s$, where the hole density achieves some value $p = p_s$, a salient point is observed. Analogous dependence $p(V_g)$ was observed earlier in Refs. [19, 20]. As seen the p_s values are significantly different for these structures: $p_s \simeq 1.3 \times 10^{11} \text{ cm}^{-2}$ for structure 1122 and $p_s \simeq 0.35 \times 10^{11} \text{ cm}^{-2}$ for structure 1023. The decrease of $|dp_H/dV_g|$ below V_g^s can result from appearance of large enough density of states at the Fermi level, which pins it. They can be the states in the secondary maxima of the valence band spectrum at $k \neq 0$ [see Fig. 1, panels (b), (d), and (e)] or some states in the barriers or at the interfaces between the layers forming the quantum well.

To understand which of these reasons is the primary one we analyze the magnetic field dependences of R_H and ρ_{xx} in classically strong magnetic fields at relatively high temperature at which the SdH oscillations are suppressed. As an example, in Fig. 4 we have presented the R_H vs B and ρ_{xx} vs B plots measured on the structure 1023 at $T = 18$ K for two gate voltages. These dependences are typical for two-type carries conductivity. The simultaneous fitting of these plots by the standard handbook expression [21] with the use of hole densities $p^{(1)}$ and $p^{(2)}$, and hole mobilities $\mu^{(1)}$ and $\mu^{(2)}$ as the fitting parameter gives reasonable agreement with the data (see Fig. 4). The hole densities $p^{(1)}$ and $p^{(2)}$ found by this way and the total density $p^{(1)} + p^{(2)}$ together with the hole Hall density $p_H = 1/eR_H(0.1 \text{ T})$ for different V_g values are plotted in Fig. 5(a). It is seen that $p^{(1)}$ is close to p_H and the $p^{(1)} + p^{(2)}$ points fall on the straight line which describes the n_H vs V_g data. The values of the mobility of the second type holes are about $(2 - 3) \times 10^3 \text{ cm}^2/\text{Vs}$ [see Fig. 5(b)], which beyond doubt correspond to a free carrier conductivity.

Thus, one can argue that the decrease of $|dp_H/dV_g|$ at

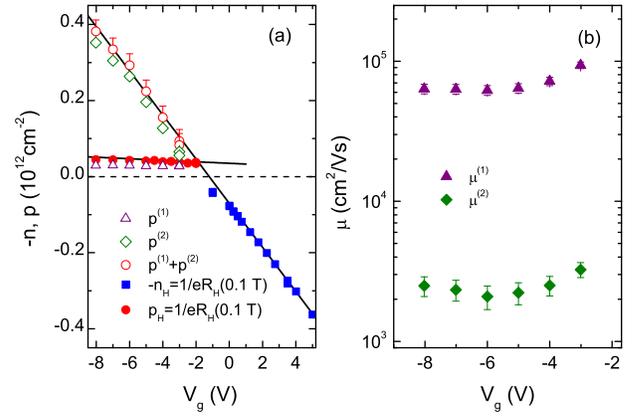


Figure 5. (Color online) The gate voltage dependences of the carrier densities (a) and mobilities (b) found from the fit of the magnetic field dependences of ρ_{xx} and R_H shown in Fig. 4.

$V_g < V_g^s$ is result of occupation of the secondary maxima of the valence band when the Fermi level approaches these maxima with the decreasing gate voltage. For this case the value of p_s should monotonically depend on the quantum well width. To compare the data obtained for the different structures with different carrier density at $V_g = 0$, different insulator thickness and, hence, the different dn/dV_g value we have plotted the values of the Hall carrier densities $1/eR_H(0.1 \text{ T})$ against the charge in the quantum well $Q = C(V_g - V_g^{\text{CNP}})$ for all the structures investigated in Fig. 6. The values p_s found from Fig. 6 are plotted against the nominal quantum well width in Fig. 7(a). The corresponding theoretical dependence calculated within framework of the 8×8 Kane model in Ref. [11] is presented also. Taking into account the uncertainty in the real quantum well width and error in p_s determination it can be concluded that the value of p_s in such a type of heterostructures can be used to estimate the quantum well width. Note, the experimental values of B_c are also in a satisfactory agreement with calculation results as evident from Fig. 7(b).

Before closing this section let us discuss the specific feature in the behavior of the hole density evident at $p > p_s$. As seen from Fig. 6, the higher is the p_s value, the larger is the slope of the dependence $p_H(Q)$ at $p_H > p_s$. Since the Hall coefficient $R_H(0.1 \text{ T})$ gives the density of holes only in the central maximum, the slope of the dependence $p_H(Q)$ at $p_H > p_s$ is determined by the relation between the densities of hole states in the central and secondary maxima ν_1 and ν_2 , respectively: $e dp_H/dQ = 1/(1 + \nu_2/\nu_1)$. Thus, the increase of dp_H/dQ with growing p_s is in qualitative agreement with the fact that the density of states in the central maximum ν_1 at the energy of secondary maximum increases with the energy increase due to the nonparabolicity of dispersion law $E(k)$.

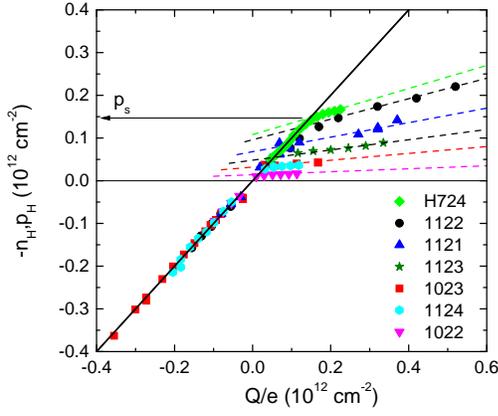


Figure 6. (Color online) The Hall carrier density plotted against the charge in the quantum well. As an example arrow shows how the p_s value is found for sample H724.

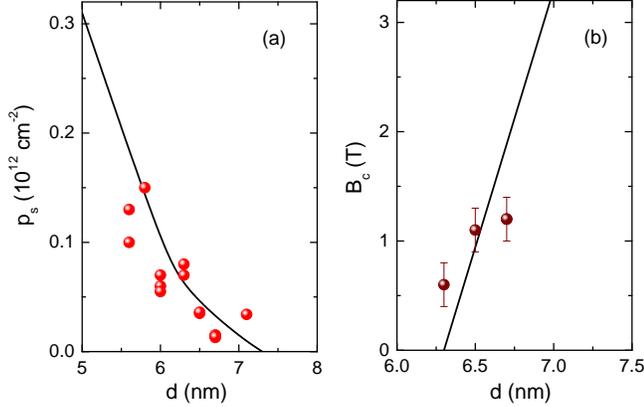


Figure 7. (Color online) The values of p_s (a) and B_c (b) plotted against the nominal quantum well width. Symbols are the data, the lines are theoretical results [11].

V. SPIN-ORBIT SPLITTING OF VALENCE AND CONDUCTION BANDS

Now let us turn to a more detailed study of the band spectra. To do this we have measured the SdH oscillations and their angle dependence over the wide carrier density range.

A. Normal band ordering

1. Valence band

Qualitatively the results obtained for structures 1122 and their interpretation are analogous to those published in Ref. [14]. Nonetheless, we briefly describe the key results that are important for the interpretation of results both for valence and conduction bands for structures with normal and inverted spectra. The gate voltage dependence of the Hall density, the dependences

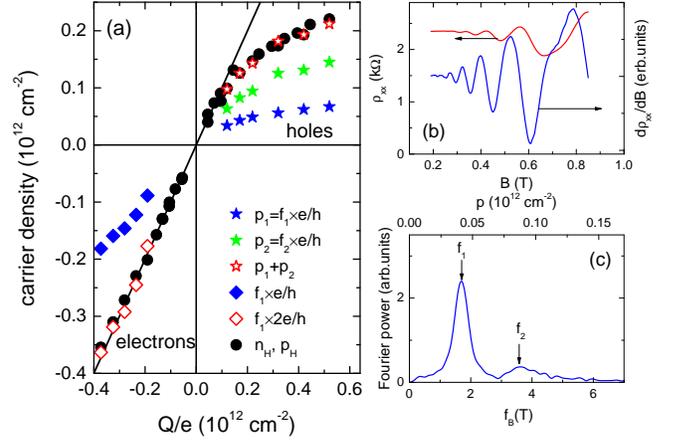


Figure 8. (Color online) (a) The carrier density plotted against the charge in the quantum well. The circles are the Hall carrier density, the stars and diamonds are the carrier densities obtained from analysis of the SdH oscillations as described in the text. (b) ρ_{xx} and $d\rho_{xx}/dB$ plotted as a function of magnetic field for $Q/e = 1.3 \times 10^{11} \text{ cm}^{-2}$. (c) The result of the Fourier transformation of the SdH oscillations shown in the panel (b). Structure 1122.

$\rho_{xx}(B)$ and $d\rho_{xx}/dB$ taken for the structure 1122 at $Q/e = 1.3 \times 10^{11} \text{ cm}^{-2}$, and the corresponding Fourier spectrum of the SdH oscillations are presented in Fig. 8.

As seen from Fig. 8(c), two maxima with the frequencies f_1 and f_2 are easily detected in the Fourier spectrum. Noteworthy is that the ratio of the frequencies is close to 2, i.e., the Fourier spectrum is similar to the case when the spin splitting of the Landau levels manifests itself with the magnetic field increase. In such a situation the carrier density should be determined as $p_{\text{SdH}} = f_1 \times 2e/h$. So, we obtain $p_{\text{SdH}} = (0.88 \pm 0.05) \times 10^{11} \text{ cm}^{-2}$ for this concrete case. However, the Hall density at this Q value is significantly larger, $p_{\text{H}} = 1.31 \times 10^{11} \text{ cm}^{-2}$ as seen from Fig. 8(a).

Such a difference between the hole density found from the SdH oscillations within the proposed model and the hole density found from the Hall effect takes place over the whole hole density range in all the structures investigated. Thus the interpretation described above does not correspond to our case.

The only interpretation that adequately describes our results is as follows. Each peak in the Fourier spectrum corresponds to the subband H1 which is strongly split by SO interaction into two subbands H1+ and H1-. In this case the “spin” degeneracy is lifted, and, hence, the hole densities should be found as $p_{1,2} = f_{1,2} \times e/h$, where indexes 1 and 2 correspond to H1+ and H1- subbands. So, the total hole density is $p_{\text{SdH}} = p_1 + p_2 = (f_1 + f_2)e/h$. The results of such a data interpretation are presented in Fig. 8(a) within whole hole density range by stars. One can see that p_{SdH} coincides with p_{H} within experimental error. Thus, we conclude that the valence band is strongly split by SO interaction so that the ratio

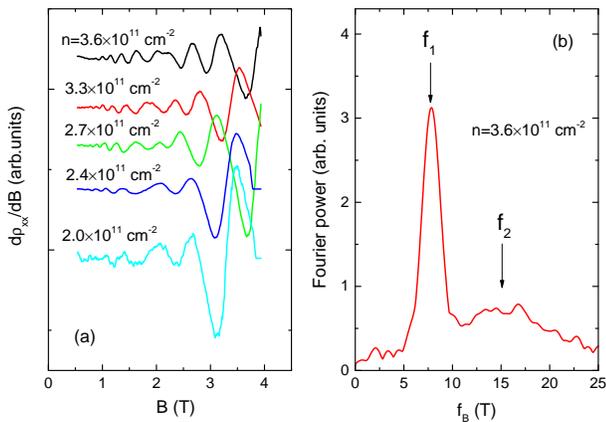


Figure 9. (Color online) The SdH oscillations of electron conductivity measured for structure 1122 at different electron densities (a) and the results of Fourier transformation of SdH oscillations for $n = 3.6 \times 10^{11} \text{ cm}^{-2}$ (b).

of hole densities in the split subbands is about 2 over the whole hole density range.

2. Conduction band

To obtain the information on the energy spectrum and splitting of the conduction band we have thoroughly studied the electron SdH oscillations [22]. In Fig. 9(a) we have presented the SdH oscillations for some Hall densities and in Fig. 9(b), as example, the Fourier spectrum for one of them. As seen from Fig. 9(b) the Fourier spectrum consists of two peaks which characteristic frequencies differ by a factor of two analogously to the case of the hole conductivity [see Fig. 8(c)]. The main contribution to the peak with the frequency f_1 comes from the low-field SdH oscillations. Therewith, unlike the oscillations in the hole domain, the electron density found as $n_{\text{SdH}} = f_1 \times 2e/h$ is very close to the Hall density as evident from Fig. 8(a) indicating that the Zeeman splitting is not resolved. The peak f_2 comes from the oscillations in the high-magnetic fields and it originates from the spin split Landau levels. This corresponds to the case when the SO splitting of the conduction band is small enough.

Thus, the analysis of SdH oscillations carried out over the whole carrier density range gives surprising result: the valence band is strongly split by SO interaction, while the splitting of the conduction band in the same structure does not reveal itself.

3. The Shubnikov-de Haas effect in tilted magnetic field

To ensure that our interpretation is correct one can explore the SdH oscillations in tilted magnetic field. Really, when the SO splitting is significantly smaller than the Zeeman energy ($\Delta_{\text{SO}} \ll g\mu_B B$), the energy of or-

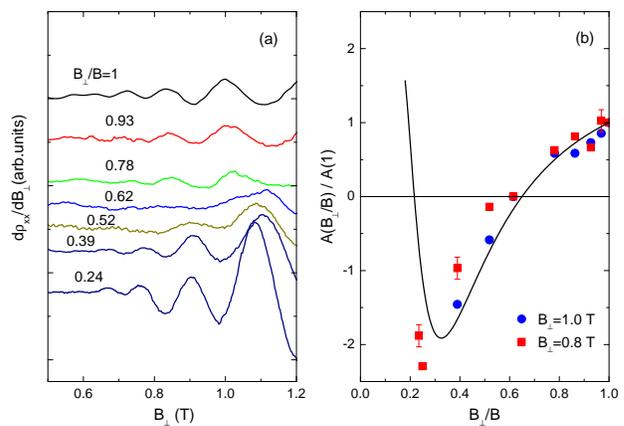


Figure 10. (Color online) The SdH oscillations of electron conductivity for structure 1122 at different B_{\perp}/B values (a) and the oscillation amplitude plotted against the B_{\perp} to B ratio for $B = 0.8 \text{ T}$ and 1.0 T (b), $n = 2.5 \times 10^{11} \text{ cm}^{-2}$. Solid line in (b) is the dependence Eq. (1) with $X = 0.32$. The inversion of the amplitude sign corresponds to the change of the oscillation phase on π .

bitual quantization depends on the normal component of the magnetic field $B_{\perp} = B \cos \theta$, where θ is the angle between the magnetic field direction and the normal to the 2D gas plane, while the energy of spin splitting depends on total B . At the low magnetic field, where the SdH oscillations are unsplit, this should manifest itself as strong angle dependence of the oscillation amplitude. In opposite case of strong spin-orbit interaction ($\Delta_{\text{SO}} \gg g\mu_B B$), the spin is rigidly coupled with the orbital motion and the SdH oscillations should be determined by the normal component of the magnetic field only.

The oscillations of $d\rho_{xx}/dB_{\perp}$ at $B_{\perp} < 1.5 \text{ T}$ for electron density $2.5 \times 10^{11} \text{ cm}^{-2}$ for several tilt angles are presented in Fig. 10(a). The oscillations within this magnetic field range are caused by unsplit Landau levels. It is clearly seen that the oscillations decrease in the amplitude when B deviates from the normal orientation, practically disappear at $B_{\perp}/B \simeq 0.6$, then change the phase and increase in the amplitude at further decrease of the B_{\perp} to B ratio. This behavior results from the change of the ratio between the spin and cyclotron energies $X = g\mu_B B/\hbar\omega_c$ with tilt angle. Supposing $g\mu_B B$ and $\hbar\omega_c$ are proportional to the total magnetic field and normal component of B , respectively, one can obtain the following expression for the angle dependence of the oscillation amplitude

$$\frac{A(B_{\perp}/B)}{A(1)} = \cos\left(\pi X \frac{B}{B_{\perp}}\right) / \cos(\pi X). \quad (1)$$

This dependence together with experimental data is shown in Fig. 10(b). One can see that a good agreement is observed when $X = 0.32$.

Let us now inspect the behaviour of the SdH oscillations with the changing tilt angle measured on the same heterostructure in the hole domain (Fig. 11). One can

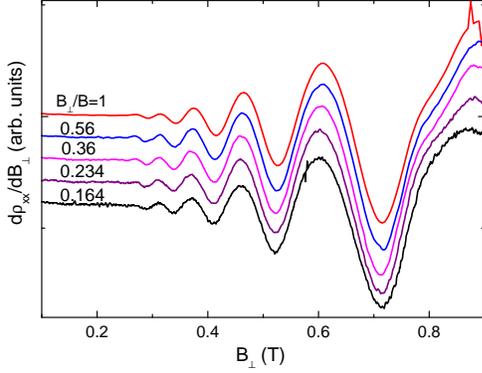


Figure 11. (Color online) The SdH oscillations of the hole conductivity plotted against the normal component of magnetic field for structure H1122 at different angles, $p = 1.35 \times 10^{11} \text{ cm}^{-2}$.

see that unlike the case of the electron conductivity, the SdH oscillations plotted against the normal component of the magnetic field remain practically unchanged with the changing tilt angle. This means that they are determined by the normal components of magnetic field only. As discussed above such a behavior should be observed when SO interaction is relatively strong: the SO splitting is much greater than the Zeeman energy. Or, alternatively, the independence of the oscillation picture on the tilt angle may be caused as well by a strong anisotropy of the effective g -factor, when the in-plane g -factor is much smaller than the perpendicular one, $g_{\parallel} \ll g_{\perp}$. However, the estimation made with the use of Eq. (35) from Ref. [23] shows that it is not the case, the Zeeman splitting differs in the perpendicular and longitudinal orientations of the magnetic field only slightly for the actual quantum well widths and actual Fermi energies.

Thus, the analysis of the angle dependences of SdH oscillations also shows that the valence band spectrum is strongly split due to spin-orbit interaction, while the conduction band remains unsplit.

B. Inverted band ordering

The dependence of carrier densities found from the Hall and Shubnikov-de Haas effects on the charge of 2D gas for structure 1023 with inverted spectrum, $d = 6.5 \text{ nm}$, is plotted in Fig. 12(a). In general, it is similar to that for structure with normal spectrum excepting the fact that the p_s value is significantly less in structures with $d > d_c$: namely $p_s \simeq 3.0 \times 10^{10} \text{ cm}^{-2}$ for structure 1023 instead of $\simeq 1 \times 10^{11} \text{ cm}^{-2}$ for structure 1122 with normal spectrum. As discussed in Section IV this results from the lesser energy distance between the top of the valence band at $k = 0$ and the secondary maxima at $k \neq 0$.

The SdH oscillations in the hole conductivity regime and their Fourier spectrum are shown in Fig. 12(b) and

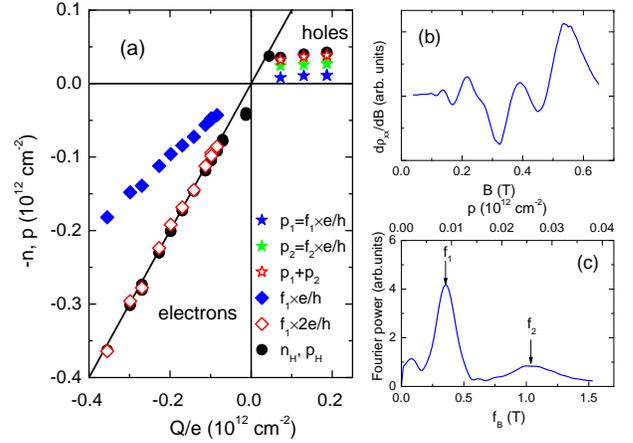


Figure 12. (Color online) Results analogous to those presented in Fig. 8, but for the structure 1023 with inverted spectrum.

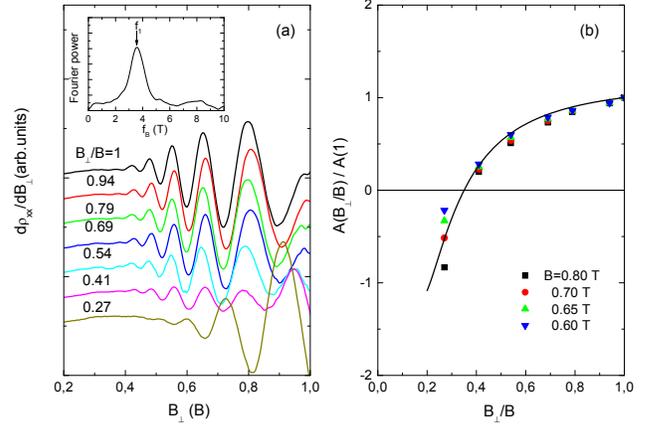


Figure 13. (Color online) The SdH oscillations of electron conductivity measured for structure 1023 at different angles (a) and the oscillation amplitude plotted against the B_{\perp} to B ratio for four B values (b), $n = 1.7 \times 10^{11} \text{ cm}^{-2}$. Solid line in (b) is the dependence Eq. (1) with $X = 0.18$. The inset in (a) is the Fourier spectrum for $B_{\perp}/B = 1$.

Fig. 12(c), respectively. As in the structures with normal spectrum, the only way, which allows us to reconcile the Hall density and data obtained from the SdH oscillations is to assume that each maximum in the Fourier spectrum is associated with the split subband.

In the electron conductivity regime, the SdH oscillations measured at different tilt angles are presented in Fig. 13(a) for the Hall density $n_H = 1.7 \times 10^{11} \text{ cm}^{-2}$. The Fourier spectrum of the oscillations at normal B orientation is shown in the inset of Fig. 13(a).

The charge dependences of the electron densities found as $f_1 \times 2e/h$ and $f_1 \times e/h$ are shown in Fig. 12(a). It is seen that $n_{\text{SdH}} = f_1 \times 2e/h$ coincides with n_H . It means that each peak of the SdH oscillations corresponds to two fold degenerated Landau level. As Fig. 12(a) shows, such a coincidence is observed over the whole range of

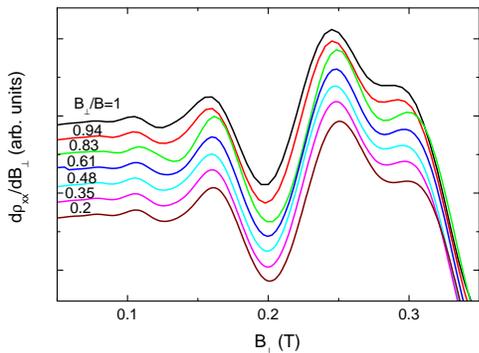


Figure 14. (Color online) The SdH oscillations of hole conductivity measured at different angles for structure H1023 with inverted spectrum, $p = 3.8 \times 10^{10} \text{ cm}^{-2}$.

the electron density. Thus, the spin-orbit splitting of the conduction band in the structures with the inverted spectrum does not reveal itself as well as in the structures with normal spectrum.

The conclusion about strong splitting of the valence band and weak splitting of the conduction band is consistent with the behavior of the oscillations in tilted magnetic field.

The behavior of the oscillations of the electron conductivity with the tilt angle shown in Fig. 13(a) is analogous to that presented in Fig. 10(a) for structure with normal spectrum. The angle dependence of the oscillation amplitude in Fig. 13(b) demonstrates that the Zeeman splitting becomes equal to half orbital one at B_{\perp}/B close to 0.35. The dependence calculated from Eq. (1) with $X = 0.18$ well describes the data [24].

The SdH oscillations of the hole conductivity measured at different tilt angle are presented in Fig. 14. As well as for the structures with normal spectrum (see Fig. 11), neither positions nor amplitudes depend on the tilt angle that corresponds to strong SO splitting of the valence band (see discussion above).

Summing up all the data given above we can formulate the main outcome. The analysis of the Shubnikov-de Haas and Hall effects over the whole carrier density range in the structures with $d < d_c$ and $d > d_c$ gives surprising result: the valence band is strongly split by SO interaction in the structures both with normal and inverted spectrum, while the splitting of the conduction band in the same structures does not reveal itself. To illustrate the written we have plotted the ratio between the carrier densities in the split subbands in the valence and conduction bands as a function of carrier density in Fig. 15 [25]. This figure is the key result of the paper.

VI. COMPARISON WITH RESULTS OF $k \cdot P$ AND TIGHT-BINDING CALCULATIONS

The strong SO splitting of the valence band at $d < d_c$ was observed in our previous paper [14]. The experi-

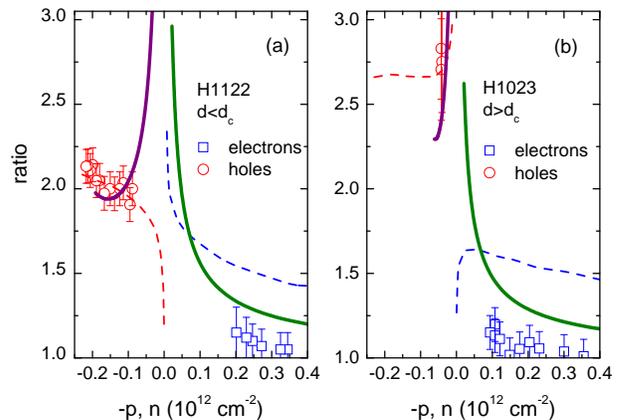


Figure 15. (Color online) The ratio between carrier densities in split subbands plotted against the carrier density for structure 1122 (a) and 1023 (b) with normal and inverted spectra, respectively. Symbols are the experimental results, the dashed line are the result of the kP calculations described in the text, the solid lines are the results obtained by the sp^3 tight-binding method [15].

mental results are interpreted under assumption that the quantum well is located in the strong electric field of a p - n junction that in its turn results in strong Bychkov-Rashba effect. It has been shown that the kP model describes the valence band spectrum quantitatively. Unfortunately, we were unable in that paper to study the conduction band with the accuracy needed to obtain the SO splitting reliably. For heterostructures studied in the present paper it was possible.

Let us now compare the experimental results obtained in the present paper both for the valence and conduction bands with the results calculated within framework of the kP model used in Ref. [14]. As before, we use here the six-band Kane Hamiltonian. The direct integration technique is applied to solve the Schrödinger equation as described in Ref. [26]. The other parameters were the same as in Refs. [27, 28]. As in the previous paper [14], the value of electric field serves as the fitting parameter, which provides the ratio between the hole densities in the split subbands, which is observed experimentally (see Fig. 15). It is approximately equal to $80 \text{ mV}/d$ for all the structures under study.

Let us now inspect the SO splitting of the conduction band. The calculated energy spectrum for the heterostructures with $d < d_c$ and $d > d_c$ is shown Fig. 16. The energy splitting Δ_{SO} as a function of the quasimomentum value is shown in the insets. It is seen that the splitting of the conduction band at low k values, $k < 0.1 \times 10^6 \text{ cm}^{-1}$, is much larger than that of the valence band for $d < d_c$. For $d > d_c$, the situation is opposite; the valence band is split much stronger than the conduction band. Such relationships between splittings agree with the known result of symmetry analysis according to which $\Delta_{so} \propto k$ for the c1 band and $\Delta_{so} \propto k^3$ for the H1 band near $k = 0$. It should be

emphasized that so low quasimomentum values correspond to very low carrier densities ($\sim 10^9 \text{ cm}^{-2}$ and less) which are inaccessible experimentally. At larger quasimomentum values, these relationships are violated. The splitting values of c1 and H1 bands become close to each other. As seen from Fig. 16 the difference in the spin-orbit splitting for the valence and conduction bands does not exceed 30 % for both heterostructures for $k_F > 0.5 \times 10^6 \text{ cm}^{-1}$ that corresponds to the actual carrier density range $n, p > 5 \times 10^{10} \text{ cm}^{-2}$ [see Fig. 15].

In order to compare the results of the kP calculations with the experimental data, we have calculated the carrier density in the split subbands as $k^2/(4\pi)$ and plotted the n_2 to n_1 and p_2 to p_1 ratios against the total carrier density in Fig. 15 by the dashed lines. It is evident that this model perfectly describes the data relating to the hole split subband and gives no agreement with the experimental results concerning the conduction band splitting. One of the reason why the used kP model fails when compared with the experiment is the fact that it ignores the asymmetry caused by the difference of the quantum well interfaces.

The role of the bulk inversion asymmetry of host crystals and interface inversion asymmetry in the SO splitting of the energy spectrum of 2D carriers was recently studied in Ref. [15]. Using the symmetry analysis and atomistic calculations the authors obtain surprising result. The asymmetry of CdTe/HgTe and HgTe/CdTe interfaces forming the CdTe/HgTe/CdTe quantum well results in the giant splitting of the energy spectrum; in the quantum wells of critical and close-to-critical width, $d \simeq d_c$, the splitting reaches a value of about 15 meV.

Analogous calculations have been performed for our concrete case of (013) HgTe. The quantum well was supposed symmetric in the sense that no electric field was applied across the the well. The microscopical strain has been calculated in the atomistic valence force field model [29] and then incorporated in the tight-binding using standard procedure [30]. The tight-binding parameters used in calculations were obtained using procedure similar to that described in details in the Supplemental Material to Ref. [15] with the following modification. In Ref. [15] it was assumed that the change of the ionicity across the interface may be extracted from the atomic levels obtained in *ab initio* calculations. Instead, one may introduce more physical parameter, the change of electrostatic potential on anion across the interface, and fit this parameter to reproduce the expected properties of the interface. Tight-binding parameters in Ref. [15] correspond to the change of electrostatic potential on Te between the quantum well and barrier equal to 1 eV, which is in accordance with *ab initio* calculations of HgTe/CdTe heteropair. For the interface between HgTe and $\text{Hg}_{1-x}\text{Cd}_x\text{Te}$ ($x \simeq 0.5$) alloy this parameter should be reduced to 500 meV which leads to proportional reduction of the interface-induced spin splitting. The carrier densities in the subbands were calculated as $S_k/(2\pi)^2$, where S_k is the area inside the Fermi contour, which is

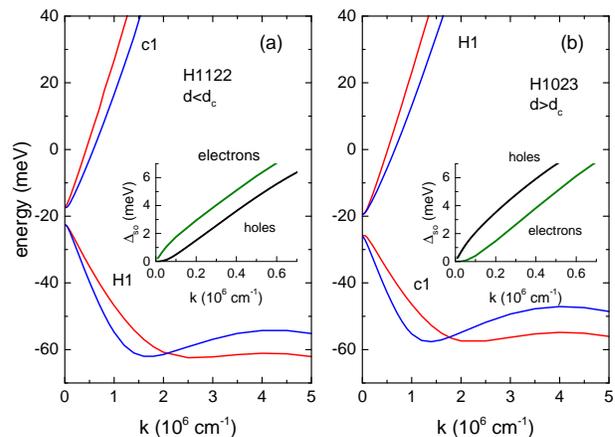


Figure 16. (Color online) The dispersion $E(k)$ for structures 1122 (a) and 1023 (b) with normal and inverted spectra, respectively, calculated within the six-band Kane model in the presence of electric field $80 \text{ mV}/d$. In the inserts, the SO splitting plotted against the quasimomentum value.

not a circle due to absence of axial symmetry. The results of calculations are depicted in Fig. 15 by the solid lines. One can see that the atomistic calculations describe the experimental SO splitting of the spectrum reasonably not only for the valence band but for the conduction band also. It is necessary to stress that no fitting parameters have been used in calculations.

VII. CONCLUSION

We have studied the energy spectrum of the conduction and valence bands of 2D states in HgTe quantum well by means of magnetotransport measurements. The structures investigated have the width of the quantum wells close to the critical value $d \simeq d_c$ where the spectrum changes from normal (at $d < d_c$) to inverted (at $d > d_c$). Simultaneous analysis of the SdH oscillations and Hall effect over the wide range of the electron and hole densities gives surprising result: the top of the valence band is strongly split by spin-orbit interaction while the splitting of the conduction band is absent, within experimental accuracy. This conclusion is supported by the results obtained in the tilted magnetic fields. The behavior of the SdH oscillations of electron conductivity with the changing tilt angle corresponds to the case when the orbital quantization is determined by normal component of the magnetic field, while the spin splitting is determined by the total field. The oscillations of the hole conductivity are totally determined by the normal component of the magnetic field indicating the SO interaction wins the Zeeman effect in the actual magnetic field range. Surprisingly, but such a ratio of the splittings is observed as for structures with normal spectrum ($d < d_c$) so for structures with inverted one ($d > d_c$). These data are inconsistent with the kP calculations which take into ac-

count the Bychkov-Rashba effect caused by electric field directed across the quantum well. It is shown that the experimental results can be reasonably described within framework of the tight-binding method which properly takes into consideration the interface inversion asymmetry.

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