

Enhanced clustering tendency of Cu-impurities with a number of oxygen vacancies in heavy carbon-loaded TiO₂ - the bulk and surface morphologies

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The over threshold carbon-loadings (~ 50 at.%) of initial TiO₂-hosts and posterior Cu-sensitization (~ 7 at.%) was made using pulsed ion-implantation technique in sequential mode with 1 hour vacuum-idle cycle between sequential stages of embedding. The final C_x-TiO₂:Cu samples were qualified using XPS wide-scan elemental analysis, core-levels and valence band mappings. The results obtained were discussed on the theoretic background employing DFT-calculations. The combined XPS-and-DFT analysis allows to establish and prove the final formula of the synthesized samples as C_x-TiO₂:[Cu⁺][Cu²⁺] for the bulk and C_x-TiO₂:[Cu⁺][Cu⁰] for thin-films. It was demonstrated the in the mode of heavy carbon-loadings the remaining majority of neutral C–C bonds (sp³-type) is dominating and only a lack of embedded carbon is fabricating the O–C=O clusters. No valence base-band width altering was established after sequential carbon-copper modification of the atomic structure of initial TiO₂-hosts except the dominating majority of Cu 3s states after Cu-sensitization. The crucial role of neutral carbon low-dimensional impurities as the precursors for the new phases growth was shown for Cu-sensitized C_x-TiO₂ intermediate-state hosts.

1. Introduction

Transition metal (TM) oxide semiconductors with a relatively narrow band-gap (2.7 eV – 3.2 eV) are considered as an ideal host-materials for the area of photovoltaic applications because of their good chemical stability, non-toxicity, low-cost and well-developed conventional synthesis methods. Among TM-oxides the titanium dioxide, having several polymorphs, seems to be potentially promising as a base for highly efficient photocatalyst in water splitting [1-2], anti-fouling [1], the major component for the photostimulated degradation of pollutants [3-5], TM-sensitization of solar-cell production [6-7]. The TM-sensitization of TiO₂-host is considered to be highly effective for the photoactivity improvement of TiO₂ by means of decreasing the charge carriers recombination rates and, hence, the enhancement of photocatalytic features. Despite the reported relatively wide commercial applicative area for TiO₂, there are still several scientific and technological challenges which are impeding the expansion and improvement of TiO₂-host functionality. Recently it was established, that an optimization attempts of TiO₂ functional properties with the use of conventional techniques lead in some cases to even a bit wider band-gap value than that for untreated host-matrix, thus indicating the incompatible with photo-energetics area the final electronic structure of the functionalized host [8-9]. For instance, Choudhury et. al. was reporting about the strong re-arrangement of initial TiO₂ atomic and electronic structure after TM-sensitization, linking this phenomenon with Jahn-Teller distortions of TiO₆-octahedra, caused by appeared oxygen vacancies (both point and hollow-types), strong *d-d* excitations and conversion from the octahedral symmetry to tetragonal one [10]. As a separate case of matter the rutile-to-anatase temperature stimulated junctions were reported and discussed as well. So it becomes clear that the functional properties of TM-sensitized TiO₂-host might be essentially affected by a large number of factors which are including the particle size (microcrystallinity type), the active surface area of the fabricated particles, the morphology of the material, etc. – viz. *in-situ* synthesis method and conditions. Wherein one from the major challenging questions is the point how to modify the energy band-gap in order to coincide well with the requirements of an "ideal" photocatalyst [11-13]

and, at the same time, not to distort the residual functionality of a modified host by suppressing the possible recombination of photoinduced electrons and holes under visible light excitation.

An applicative importance of titanium dioxide as a host-material for the structure sensitization in terms of photoactivity is following from the fact that this material is allowing to employ the relatively wide range of processing techniques with promising final results. For instance, an electrochemical synthesis (an anodizing of titanium in different solutions and doping with transition metals [14]) seems to be promising because of the yielded final functionalized TiO₂-structure with enhanced photocurrent [15-16], deceleration of the photocatalytic degradation side-effects [17] and, thus, elongation of performance cycle-time, in some cases of corrosion protection [18] of an active photocatalyst electrode zones. With that, this materials science area seems to us as still challenging because of the need to apply actually specific solvents [16,19]; appearance of oxygen sublattice defects in the sensitized TiO₂ [19]; the unwanted variety of the charge states of an employed for the sensitization transition-metal dopant [20].

As a relatively novel and effective approach for TiO₂ sensitization with TM's, an additional doping with carbon might be considered [21-27]. The most remarkable conclusions, following from these cited papers allow to point the reduction of TiO₂ microparticles dimension with simultaneous enlarging of their photoactive surface (the so-called synergistic effect, see i.e. Refs. [28-30]), higher absorptive ability of the incident light and the re-distribution of photo-generated charge-carriers with the overall improvement of energy conversion efficiency. Unfortunately an excess of embedded carbon via high-concentration mode of carbon-loadings of host-matrix might introduce concrete and purely negative side-effect: the increasing the overall electric resistivity of a carbon-doped photoactive material with the onward degradation of photoinduced current up to negligibly small value [31]. In order to avoid the overall electric resistivity increasing of C-TiO₂ host, the co-doping with a good TM-conductor might be considered, i.e. metal copper, but this will increase of amorphous grain boundaries with all the consequences (at least while employing conventional doping approaches) [10]. Thus the

situation with carbon-copper co-doping of TiO₂-host is not so straightforward and the combined theoretical and experimental study of this question is needed for the deeper clearance.

In the current paper we will present the examined and established influence of heavy carbon-loadings (~ 50 at.%) and posterior Cu-sensitization (~ 7 at.%) on the electronic structure and defects of initial *bulk* and *thin-film* (surface) morphologies of TiO₂-hosts. As previously, the combined X-ray Photoelectron Spectroscopy (core-levels and Valence Band) and DFT-based modeling approach will be employed. The key-point attention will be paid to the clustering tendency of embedded carbon and copper implants as well as to the oxygen sublattice re-arrangements. This pending case of TiO₂ carbon-loadings, being the characteristic feature of applied ion-implantation techniques, essentially exceeds the standard chemical solubility limits of conventional carbon-doping modes (over the threshold saturation) for titanium dioxide and, thus, might be of special interest for different technological applications.

2. Samples preparation and experimental details

For the subsequent ion-implantation treatment of TiO₂ as a host-matrix, the *bulk ceramics* and *thin-film* morphologies of titanium dioxide were employed as previously. Recall, that using sample the preparation methodology, XRD, and preliminary XPS characterizations, reported in Refs. [32-34], the single-phase phase rutile with lattice parameters of $a = 4.592\text{\AA}$ and $c = 2.960\text{\AA}$ and average crystallite size of about 200 nm (*bulk ceramics* or from now simply *bulk*) and single-phase anatase *thin-films* [32] with lattice parameters of $a = 3.785\text{\AA}$ and $c = 9.487\text{\AA}$ were obtained.

The exordial brief analysis of third-party research data allows to derive that the carbon-copper co-implantation treatment seems to be the logically justified choice for employed TiO₂-hosts. Since the simultaneous pulsed co-implantation with two types of ions adducts to the essential conversion of an oxygen-ligand and central atom sub-lattices (see i.e. Ref. [35]), thus one cannot assume that this mode of implantation will be suitable for the electronic structure soft modulation (ESM) [36] of our TiO₂-hosts [37]. So we will make an attempt to perform the sequential step-by-step pulsed ion-implantation firstly with carbon (so-called C_x→TiO₂ loadings [38]), but without hydrophobic

treatment of the surface which usually causes the appearance of "alien" OH⁻ groups and then with copper. The main idea of such an approach is to prevent or, at least, to reduce the initial host-structure total conversion for better compatibility with technological ESM-conditions.

The sequential carbon (first step of separate embedding) and then copper (second separate step) ion-implantations of TiO₂-hosts in a pulsed-repetitive mode per each step were made at a vacuum pressure of 3×10^{-3} Pa. The MEVVA-type ion-source provided the identical embedding conditions both for C-ions and, after that, for Cu-ions: ion-energy of 30 keV, the pulsed-beam current density of 0.65 mA/cm², pulse repetition rate of 12 Hz, a pulse duration of 0.3 ms and flux of 1×10^{17} cm⁻². Total time-exposure under appropriate ion-beam treatment was not less than 5 hours for primary carbon embedding and then 70 minutes for the onward copper embedding with 1 hour of an idle-cycle between carbon and copper loadings. The average temperature of the host-samples under ion-embedding did not exceed 300°C. The earliest and detailed report about principles of the above briefly described technique and the current state of matter are presented in Refs. [39-40], respectively.

X-ray photoelectron qualification of the final electronic structure of sequentially carbon-copper ion-beams treated TiO₂-hosts was performed using a PHI XPS Versaprobe 5000TM spectrometer (ULVAC-Physical Electronics, USA) based on a classic X-ray optic scheme with a quartz monochromator and hemispherical energy analyzer working in the range of binding energies from 0 to 1500 eV with an employment of Al K α radiation (1486.6 eV, $\Delta E_{\text{resolution}} \leq 0.5$ eV, 100 μm X-ray probe dia, X-ray power loading of the samples less than 70 Watts). More detailed description of the used XPS recording parameters might be found in our previous publications (see, i.e. Refs. [33-35]). For a complete description of PHI XPS Versaprobe 5000TM system, please, see the web-site of spectrometer manufacturer. According to the accepted ASTM Standards of XPS measurements and charge-referencing [41] the complete electronic structure mapping – fast wide-scan (survey), core-levels and valence bands – had been made. The interpretation of the XPS data obtained was performed with the crossed-check referencing of internationally accepted XPS Databases [42-43] and Ref. [44]. No argon sputtering of the samples under study was made to prevent an inadvertently sputter-induced

fabrication of metal carbides and carbonates which might occur under this kind of treatment in carbon-loaded metal-containing oxides, nitrides, borides, etc. [43]. As a separate and additional testimony that no TiC or/and other titanium-carbon derivative compounds are present in our host-samples after carbon-ion loadings an exactly high binding energy difference between the Ti 2p XPS for TiO₂ and TiC can be considered – about 5.3 eV [45]. This noticeable difference will be clearly visible even in the XPS wide-scan spectra but it is absent, and all Ti 2p for our implanted hosts and reference TiO₂ have the same BE-positions within the spectral resolution of recorded XPS wide-scans (see Fig. 1).

The comparison of XPS fast wide-scan qualification (XPS chemical elements analysis) for sequentially carbon-copper ion-beams modulated TiO₂-hosts in the *bulk*, and *thin-film* morphologies are shown in Fig. 1. One can see that the XPS wide-scans of the samples under study are

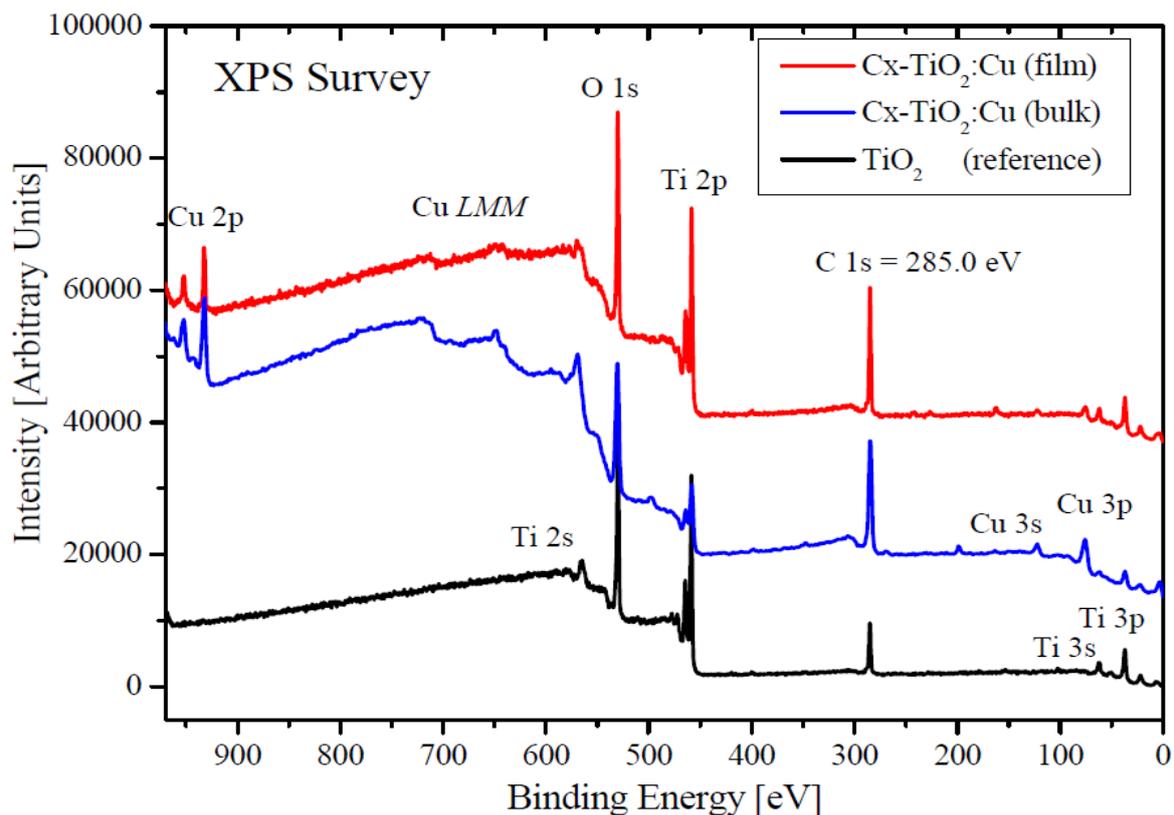


Figure 1. XPS fast wide-scan (survey) elements analysis of the carbon-copper modulated TiO₂-hosts in the *bulk* and *thin-film* morphologies with referencing to TiO₂ XPS external standard.

characterized by the appearance of additional specific XPS peaks, belonging to the appropriate copper electronic states. Moreover, much higher C 1s signal is manifesting in compliance with TiO₂ XPS

external standard. The detected spectral hangs are not abrupt because exactly these elements had been sequentially implanted into TiO₂-hosts. Leastways, around three times higher intensity of C 1s in the wide-scans of C_x-TiO₂:Cu samples in both employed morphologies versus reference TiO₂ might be recognized as a sequence of performed carbon loading via first ion-implantation step. Nevertheless, there is a chance to speculate reasonably about vacuum hydrocarbons contamination of our samples which usually results in distortions of XPS C 1s simple-shape symmetry as well as emerging additional well-known XPS sub-bands in the range of C 1s binding energies. To clarify this case of matter the high-resolution C 1s core-level XPS analysis was made (see Fig. 2).

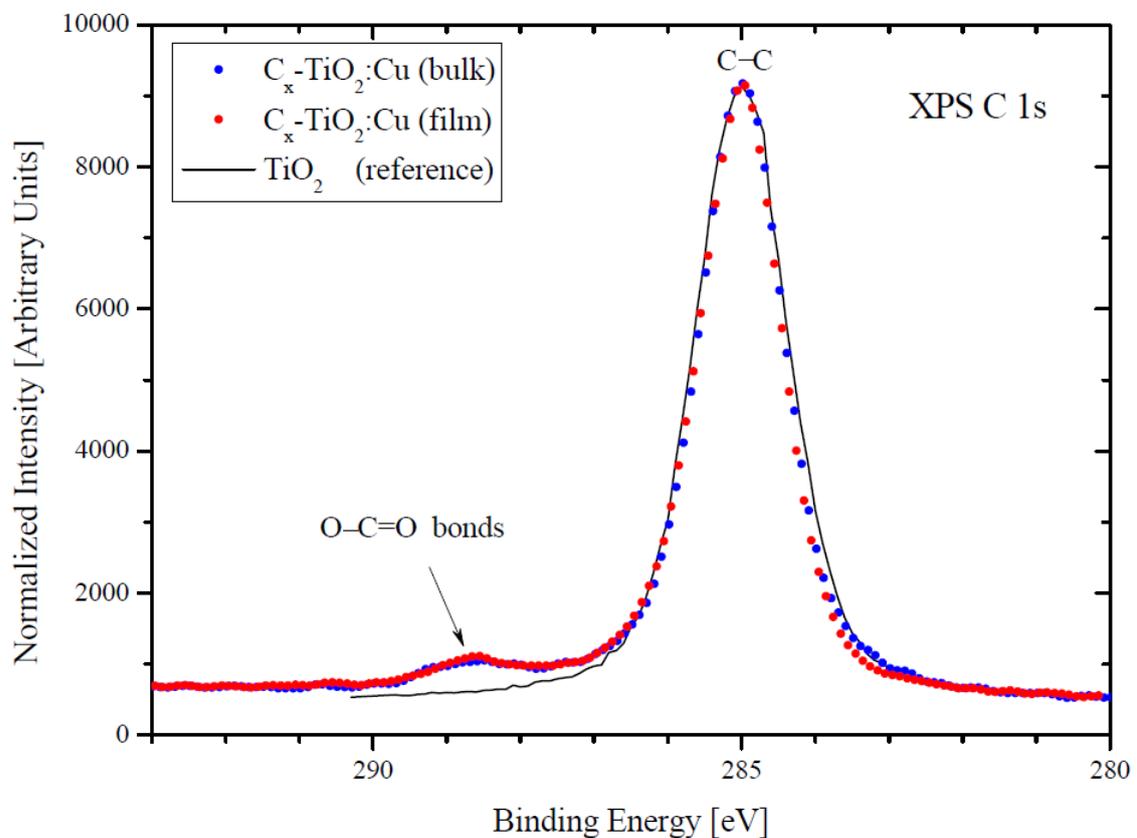


Figure 2. XPS C 1s core-level spectra of the carbon-copper modulated TiO₂-hosts in the *bulk* and *thin-film* morphologies with the referencing to TiO₂ XPS external standard.

From Figure 2 one can clearly see, that C 1s core-levels in the electronic structure of C_x-TiO₂:Cu samples and TiO₂ reference have the identical in shape-symmetry XPS main-lines located at 285 eV which represent the C–C bonding. With that, there are no any additional spectral contributions in the

vicinity of 286 eV binding energy value, thus no chance for the statements about C–O–C type of contaminations or/and imperfections, because they are manifesting themselves, in general, as a visually asymmetrical shoulder exactly in this BE-region for the wide set of carbon-oxygen containing materials [42-43]. At the same time, the 288.7 eV XPS sub-band is present only in the carbon-copper modulated $C_x\text{-TiO}_2\text{:Cu}$ samples with no even a sign of it in the C 1s core-level of TiO_2 -reference (Fig. 2). Thus the character of 288.7 eV XPS sub-band manifestation, in our humble opinion, might be linked with the aftermath of carbon loadings via ion-implantation and not with vacuum hydrocarbons contamination. According to XPS Databases, this sub-band is arising due to the irregular O–C=O fragments [43] thus supporting our derivation about C-implanting reason for it. Recall that the 288.7 eV XPS sub-band is absent at all in C 1s spectrum of TiO_2 reference, which was not carbon-implanted, but only surface-marked with adventitious carbon for the proper charge referencing (see Fig. 2). From above one might conclude rather about an incomplete interaction of implantation- embedded carbon with titanium-oxygen lattice than about vacuum hydrocarbon contaminants, as it was reported in Ref. [42]. Again, no any XPS signals were found in the range of 282 eV in C 1s core-levels of the samples under study, so this is one more evidence that no TiC and/or other titanium-carbon derivatives were fabricated unexpectedly in our samples because of the improper and unsuitable technological treatments [45].

Finally, the performed XPS qualification allows to state that no alien XPS components are present in XPS wide-scans and C 1s spectra, thus proving the absence of contaminants in $C_x\text{-TiO}_2\text{:Cu}$ final samples. The averagely estimated concentrations of injected via implantation carbon and copper are shown in Table I.

Table I. Estimated concentrations of components in sequentially carbon-copper implanted TiO_2 -hosts.

Sample (morphology)	Concentration (at. %)			
	C	O	Cu	Ti
$C_x\text{-TiO}_2\text{:Cu}$ (bulk)	45.6	32.8	6.9	14.7
$C_x\text{-TiO}_2\text{:Cu}$ (film)	43.3	33.8	6.7	16.2

3. Computational method and models

Density-functional theory (DFT) calculations were performed using the SIESTA pseudopotential code [46] as had been used successfully for related studies of impurities in the *bulk* and *thin-film* morphologies of TiO₂ [40,41]. All calculations had been made employing the Perdew-Burke-Ernzerhof variant of the generalized gradient approximation (GGA-PBE) [47] for the exchange-correlation potential accounting the Dipole Correction [48] as well. After that the calculated atomic positions were completely optimized. The ground electronic state was consistently found during optimization using norm-conserving pseudopotentials [49] for the cores and a double- ξ plus polarization basis of localized orbitals for C, Cu, Ti, and O. The forces and total energies were optimized with an accuracy of 0.04 eV/Å and 1.0 meV, respectively. All calculations were carried out with an energy mesh cut-off of 300 Ry and a k -point mesh of 6×6×4 in the Monkhorst-Pack scheme for the *bulk* and 6×6×2 for the "surface" (*thin-film*) [50]. In order to plot the final Densities of States (DOSes), the k -point mesh was increased up to 8×8×6 and 8×8×4, respectively.

The employed for the model calculations supercells were including 96 atoms both for the *bulk* rutile-TiO₂ and for the *thin-film* (surface) anatase-TiO₂, but in the latter case the supercell was used as a slab because it is the most physically feasible model of (001) TiO₂ surface [33,34]. Additionally, for the *bulk* morphology used in experiment both rutile and anatase polymorphic structures were theoretically tested and carefully examined, whereas for *thin-film* (surface) the anatase polymorphic structure was selected solely (see the reason in experimental part above).

The calculations of the formation energies (E_{form}) were based on the following formula:

$$E_{\text{form}} = [E_{\text{total}} - (E_{\text{matrix}} - nE_{\text{removed}} + mE_{\text{added}})]/m \quad (1)$$

where E_{total} – is the total energy of the system with the current configuration of defects; E_{matrix} – denotes the total energy of the system before substitution of host-atoms with impurities, or the fabrication of vacancy or/and interstitial defects (the conversion mode of structural configuration), E_{removed} and E_{added} – are the energies of removed and added atoms which were calculated for the bulk

(graphite state for carbon) and gaseous phases (triplet state for oxygen) in the ground configuration; n and m – are the numbers of removed and added atoms, respectively. We have to note, that for the structures with multiple charge-state defects (i.e. embedded Cu-impurities and simultaneous presence of oxygen vacancies), E_{matrix} already denotes the energy of TiO_2 supercell with oxygen vacancies.

Based on our previous theoretical experience (see Refs. [33,34]), we decided to consider as possible copper impurities the following types of defects: substitutional (S), interstitial (I) and their combination of (xS+I) type. The case of multiple defects combination was analysed and examined separately, employing different distances between various defects – i.e. among the oxygen vacancy, carbon and copper impurities). The calculations of the formation energies were performed only for the configurations with the lowest total energies of the systems.

4. Results and Discussion

From above performed XPS qualification of C 1s core-levels in the electronic structure of $\text{C}_x\text{-TiO}_2\text{:Cu}$ samples in both studied morphologies becomes clear, that the majority contribution of carbon into our hosts turns out in the sp^3 -bonded form [50]. This means that there were weak and insufficient interactions of carbon with titanium-oxygen lattice which results, with the high probability, only in the fabrication of O–C=O clusters, and in general, an embedded carbon remains in the form of neutral C–C bonds. No any spectral signatures of carbon-titanium interactions were found (see discussion above). The XPS asymmetrical and weak contribution of O–C=O XPS bands into C 1s core-levels at 288.7 eV (see Fig. 2) allows to suppose about the minority holdings among oxygen-carbon with sp^2 -character of the latter, because all sp^2 -bonded carbon compounds have definitely asymmetrical XPS line-shapes of C 1s core-levels in compliance with high-symmetrical XPS line-shapes for sp^3 -bonding [43]. Thus from established dominating neutrality of embedded carbon logically follows, that the majority of the "host oxygen – implanted element" interactions might be accounted for the sequentially embedded copper, so we will analyze now the XPS Cu 2p core-levels of our samples which are shown in Fig. 3.

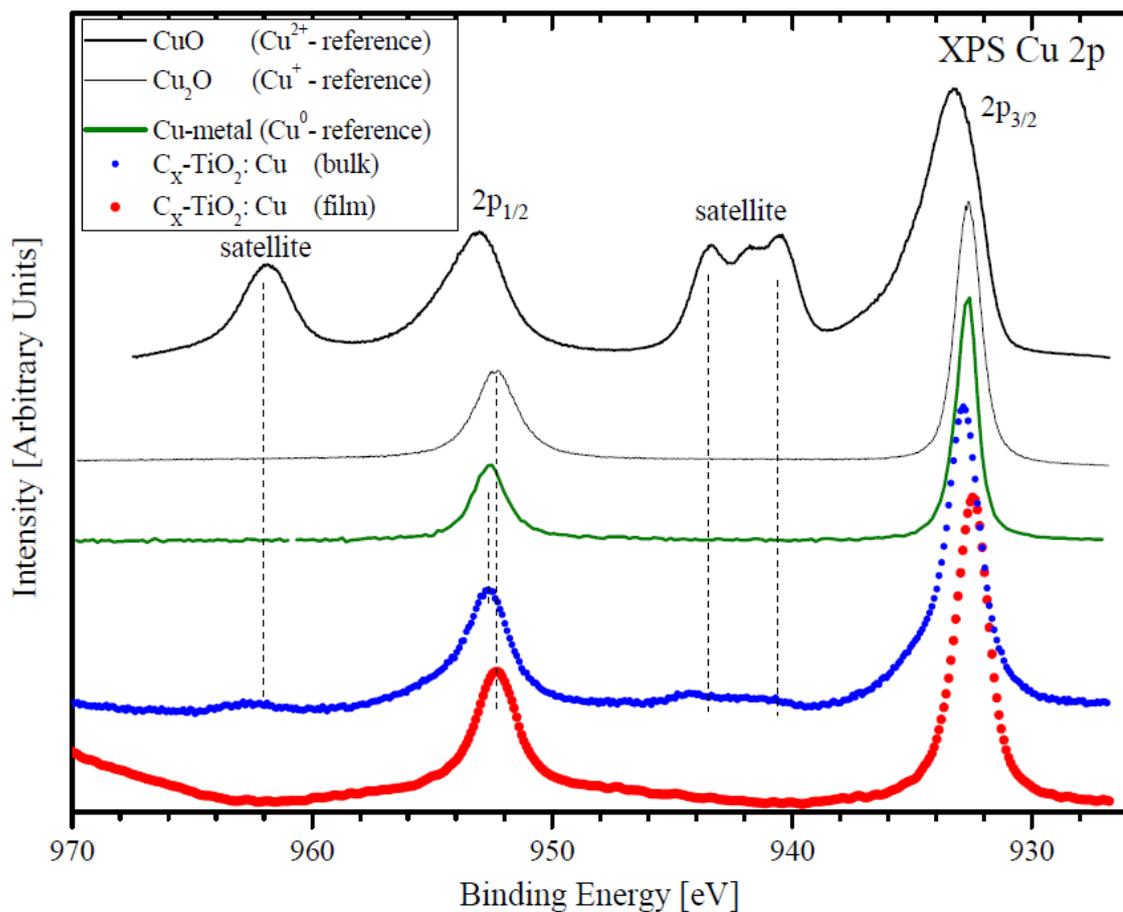


Figure 3. XPS Cu 2p core-level spectra of the carbon-copper modulated TiO₂-hosts in the *bulk* and *thin-film* morphologies with the appropriate XPS external standards: CuO, Cu₂O and Cu-metal.

One can see from Figure 2, that the main difference in XPS Cu 2p_{3/2-1/2} for C_x-TiO₂:Cu in both employed morphologies is manifested in the clearly visible BE-shift of 2p_{3/2-1/2} components relatively to each other and weak-intensity broad spectral features located in the 941-943 eV range and at 962 eV, respectively. As another visually observable feature of the *bulk* C_x-TiO₂:Cu XPS Cu 2p_{3/2-1/2} might be considered the strongest asymmetry which is causing the shoulder at 934.7 eV, if comparing with that for *thin-film* (see Fig. 3, blue and red spectra). It seems that these mentioned above well-recognizable XPS spectral features are specific to establish the valence states of embedded copper ions into C_x-TiO₂ matrices. Formally we are dealing now exactly with C_x-TiO₂ hosts because the implantation was

sequential and primarily the carbon was injected into initial TiO₂-hosts. On the strength of main-lines BE-positions in the XPS Cu 2p of our samples (Cu 2p_{3/2} = 932.9 eV for the *bulk* and Cu 2p_{3/2} = 933.4 eV for the *thin-film*) there are no background to link our XPS data with that for "copper-carbon" and "copper oxide-carbon" compounds (compare, i.e. in Cu²⁺-carbonate di-hydroxide Cu 2p_{3/2} = 934.7eV [43], in nanostructured copper phthalocyanine-sensitized multiwall carbon nanotube films Cu 2p_{3/2} = 935.9 eV [51] and in composite copper oxide/chitosan/carbon fibers Cu 2p_{3/2} = 935.7 eV [52]). Additionally, the shape of the characteristic "shake-up" satellites in these compounds is essentially dissimilar [43,52] with what was established in our XPS Cu 2p core-level analysis for both morphologies of C_x-TiO₂:Cu. Thus we have all the reasons to believe that no "copper-carbon" and "copper oxide-carbon" compounds were fabricated after sequential Cu-implantation of C_x-TiO₂ hosts on the second stage of ion-beam stimulated synthesis. Already well-known details about the XPS "shake-up" satellite and mixed valence states of copper in copper containing compounds are reported, for example, in Refs. [42-44, 53] or elsewhere. Finally, the high-oxidation and high-acidification abilities of copper (as well as of Pb, Sn, ..., etc.) are excluding any residual doubts that copper-oxygen interaction is much easy than copper-carbon one, at least without any specially applied thermodynamic conditions which were absent in our synthesis.

From above discussion preamble the oxidation roots of Cu 2p core-levels shape-transformation in the electronic structure in C_x-TiO₂:Cu samples might be supposed as the basic background. Indeed, the XPS Cu 2p_{3/2-1/2} for *thin-film* (Fig. 3, red spectrum) has the same BE-locations and line-shapes as Cu₂O and the absence of characteristic "shake-up" satellite, so formally the valence states of copper might be accounted as Cu⁺. At the same time a little larger FWHM (more at ~ 0.3 eV) of the main XPS peaks in *thin-film* C_x-TiO₂:[Cu⁺] than in Cu₂O possibly occurs due to partial contribution of metallic copper which has essentially neighbouring to Cu₂O binding energy values, making quite difficult to analyze the real state of matter in terms of copper charge states detection. This partial metallic Cu-loss effect has, at least, two valuable reasons in our humble opinion: (i) some oxygen-ligands from the titanium-oxygen lattice of the initial TiO₂-host had been previously "captured" by sp²-bonded carbon;

(ii) an employed $1 \times 10^{17} \text{ cm}^{-2}$ ion-flux for copper embedding is the ride-line fork for the ion-beam stimulated processes of small metallic clusters growth and Oswald ripening (also known as particles coarsening) [54]. Both oxygen deficit after the first step of sequential carbon implantation and a relatively low atomic fraction of copper (see Table I) are reasonable enough to cause the appearance of metallic copper small clusters on the "growth stage" according to Ref. [54]. Thus they are contributing to the XPS Cu $2p_{3/2-1/2}$ core-level of *thin-film* as well, and we suppose the final formal formula as $\text{C}_x\text{-TiO}_2\text{:}[\text{Cu}^+][\text{Cu}^0]$.

Another situation occurs with the *bulk* sample. Here one can see all the XPS signatures of the presence of Cu^{2+} electronic states. They are: (i) visually noticeable asymmetry of Cu $2p_{3/2}$ main-line (usually absent in stoichiometric Cu^+ -compounds); and (ii) the appearance of characteristic "shake-up" satellite (similar to CuO) in the same BE-region. Nevertheless, the main XPS lines for this sample are narrower than in ordinary (or natural) CuO, and the satellite is very weak (see Fig. 3). Moreover, the binding energies of the main lines are in better agreement with that for CuO_2 XPS external standard (compare thick black and blue spectra in Fig. 3), so this might be a sign of a mixed valence states of copper in the *bulk* morphology of $\text{C}_x\text{-TiO}_2\text{:Cu}$, namely $\text{C}_x\text{-TiO}_2\text{:}[\text{Cu}^+][\text{Cu}^{2+}]$. This is very close to that had been previously established by us and reported in Ref. [55]. The main dissimilarity of the XPS data reported in the current paper with that in Ref.[55] (in brief: spherical shock-waves loading by means of high-energy He^+ ions bombardment and shear under high-pressures of a native CuO as a model object) is that we got in the *bulk* $\text{C}_x\text{-TiO}_2\text{:Cu}$ the very beginning of Cu^{2+} valence states fabrication under low-energy sequential pulsed ion-beams treatment (ESM-mode) and the Cu^+ -containing phase is dominating in the *bulk* sample. Finally, our XPS data on Cu $2p_{3/2-1/2}$ for $\text{C}_x\text{-TiO}_2\text{:Cu}$ in both morphologies well agrees with that have been reported by Y.Xu *et. al.* [56] for solely Cu-doped TiO_2 , thus proving our current and above reported findings that carbon is not affecting so much the charge states of copper in TiO_2 matrix and the valence states of Cu are determined exactly by the morphology of TiO_2 matrix and its oxygen ability to interact with embedded copper in our case of implantation. This conclusion does not seem precarious, because the same C-modulation conditions

were applied both for TiO_2 *bulk* and *thin-film* hosts and, as a host reply result, the same reason for the manifestation of 288.7 eV XPS sub-band in these matrices occurs. With that, the morphology (*bulk* versus *thin-film*) and polymorphism (anatase versus rutile) of employed matrices are different because of the dissimilar titanium-oxygen unit-cell parameters in these polymorphs (slightly distorted cubic close-packing for O with Ti in a half of O_h holes in anatase versus distorted hexagonal close-packing O with Ti in a half of O_h holes in rutile [57]). The latter might be at least one of the possible valuable reasons for different type of Cu-embedding into the polymorphic hosts with the identical chemical-composition formula. This matrix-reply effect for the metal-implant with easily oxidized abilities was established by us as well for the Pb-modulated TiO_2 -hosts being in the same morphologies [53] and does not contradict the current findings.

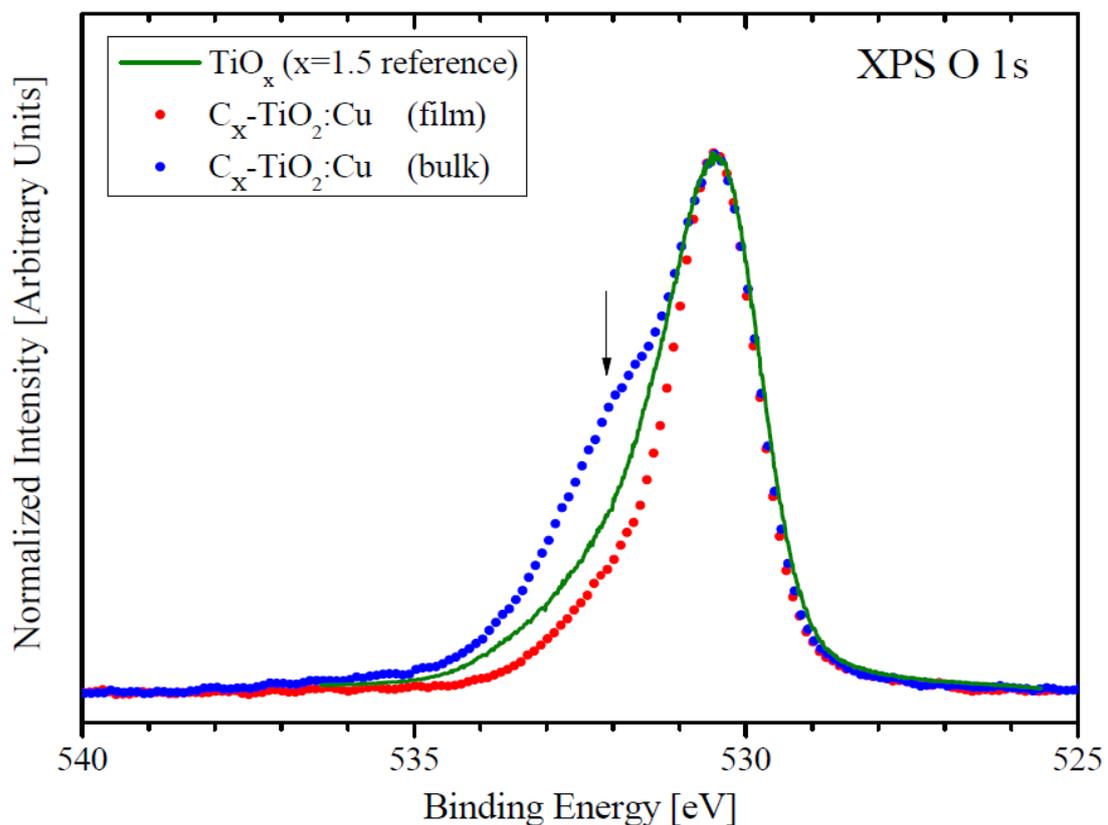


Figure 4. XPS O 1s core-level spectra of the carbon-copper modulated TiO_2 -hosts in the *bulk* and *thin-film* morphologies and TiO_x ($x = 1.5$) XPS external standard.

The recorded XPS O 1s core-level spectra of the samples under study and TiO_x ($x = 1.5$) reference are shown in Fig. 4. The actual BE-locations for these spectra were established as the following: $\text{TiO}_x = 530.4$ eV, *bulk* $\text{Cx-TiO}_2\text{:Cu} = 530.1$ eV and *thin-film* $\text{Cx-TiO}_2\text{:Cu} = 530.3$ eV. In order to analyze precisely the O 1s core-level line-shapes deviations from each other, we had shifted them in BE-scale employing the manner of the best match with TiO_x ($x = 1.5$) XPS O 1s external standard, as it had been suggested for the first time as a methodical approach in XPS analyzing by J. Kawai *et. al.* [53]. The use of TiO_2 O1s XPS external standard for comparing seems unreasonable, because carbon-embedding before copper implantation stage had already produced some oxygen defects in the microstructure of matrices (ballistic effects of an ion-implantation), and, thus, it is more correctly to speak about Cu-implantation into Cx-TiO_x matrices, rather than into Cx-TiO_2 . As an additional valid argument for our XPS external standard selection the 529-530 eV BE-range for stoichiometric metal oxides [43-44] might be considered, whereas our XPS O 1s data is not matching it. Concretely, the stoichiometric TiO_2 XPS external standard gives the value of 529.8 eV for O 1s core-level, which is, from the very beginning, actually away from the neighboring BE-positions for the TiO_x ($x = 1.5$) and $\text{Cx-TiO}_2\text{:Cu}$ samples in both used morphologies so we haven't displayed it for the Figure 4 clarity.

From Figure 4 one can see the dissimilar behavior of known oxygen defects-linked 532 eV XPS sub-band. Recall, that in our case there is no background to interpret this sub-band on the basis of OH-groups, because no hydrophobic treatment of the surface of our samples which usually causes the appearance of "alien" OH groups was used (see samples synthesis descriptions in "Samples preparation and experimental details" section). On the oxygen defects-linked basis, the *bulk* $\text{Cx-TiO}_2\text{:Cu}$ has the strongest deviations from initial O – Ti bonding in untreated TiO_2 and even from oxygen-deficit TiO_x ($x = 1.5$), whereas *thin-film* $\text{Cx-TiO}_2\text{:Cu}$ has the smallest one (compare blue and red O 1s spectra at Fig. 4). This established strong XPS dissimilarity of O 1s core-level spectra for the *bulk* and *thin-film* morphologies well agrees with the above reported XPS Cu $2p_{3/2-1/2}$ core-level analysis, according to which the final formal formulas are accounted as *bulk* $\text{Cx-TiO}_2\text{:} [\text{Cu}^+][\text{Cu}^{2+}]$ and

thin-film Cx-TiO₂:[Cu⁺][Cu⁰] and that's why. The presence both of Cu⁺ and Cu²⁺ in the *bulk* Cx-TiO₂ means that all embedded copper ions "captured" the oxygen from TiO_x oxygen-titanium lattice, increasing the oxygen deficiency in already oxygen deficient Cx-TiO₂:Cu up to the highest possible level (and, thus, the 532 eV sub-band has the most profiled XPS parameters in O 1s core-level). Conversely, in the *thin-film* only Cu⁺ states are linked with oxygen, because Cu⁰ evidently can't be bonded with it at all, so the 532 eV defective sub-band has the lowest XPS profiling (see Fig. 4). With that, the overall line-shape of XPS O 1s spectra exhibits the TiO_x-character majority (i.e. oxygen-deficient character) and is quite dissimilar with that for stoichiometric copper oxides. Low concentrations of embedded copper are responsible for it (see the data in Table I).

On the first look the appearance of Cu⁰ means Cu-losses from the ion-beam stimulated synthesis and modulation of Cx-TiO₂ microstructure (this effect was established earlier by the other researchers even for amorphous hosts [59,60] and in our previous findings [61,62]), but, what might be really interesting in our humble opinion, it is accompanied by embedded metal treating of oxygen vacancies (see, for instance, Ref.[63,64]), and this occurred one more time with the current samples even under sequential and relatively soft (regarding applied mode) ion-beam stimulated synthesis. Thus one might suppose about more deep and strong coupling among these two effects, because they were detected and fixed in appearing for a relatively wide set of matrices and implants of a dissimilar chemical origin, but, at the same, are pooled by the common type of nonequilibrium ion-embedding conversion and functionalization of the final microstructure of a material. Lastly, we have to note as well that it was not possible to record the XPS O 1s spectra of Cx-TiO₂ hosts in both morphologies before Cu-implantation step due to technical limitations of the sequential carbon-copper loadings (uninterruptable idle-cycle in vacuum before Cu-implantation without a contact with air ambient in order to prevent possible accidental additional and undesirable oxidations or/and acidifications of Cx-TiO₂).

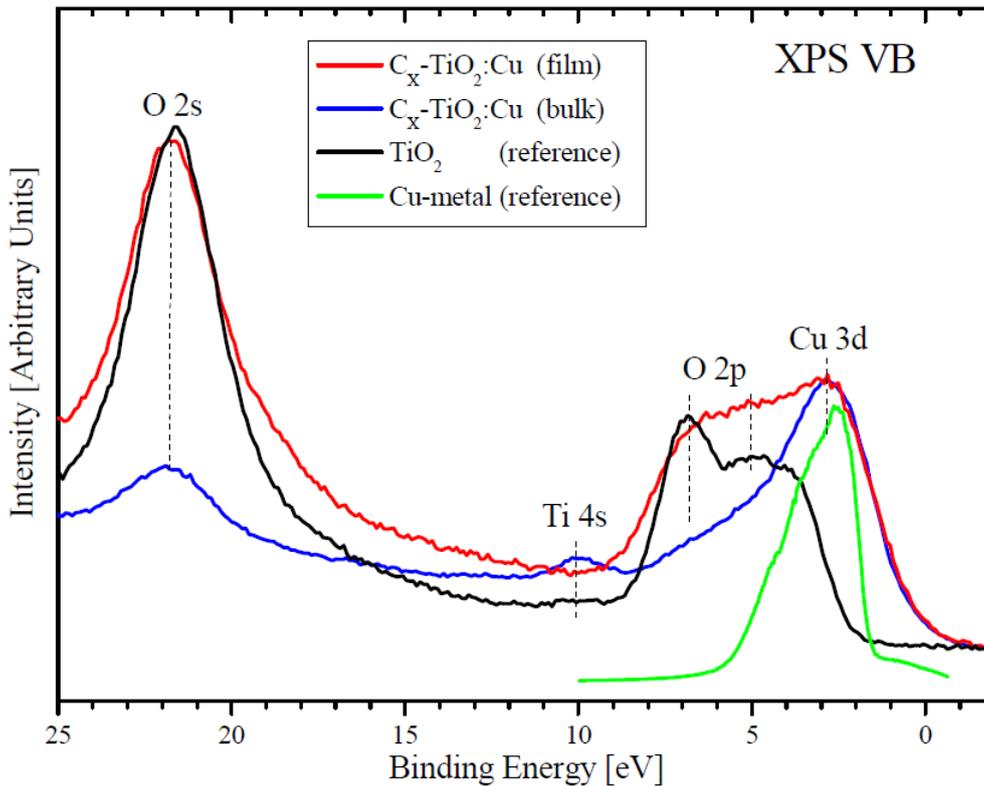


Figure 5. XPS Valence Band (VB) spectra of the sequentially carbon-copper modulated TiO₂-hosts in the *bulk* and *thin-film* morphologies, TiO₂ and Cu-metal [65-66] XPS external standards.

Valence band XPS spectra of the samples under study and XPS external standard references are shown in Fig. 5. On the whole, there is no at all any BE-shift of O 2s core-like (or semi-core) partial electronic states and the formal differences in O 2s band-tails are rather linked with dissimilar XPS BackGround (BG) contributions because of the dissimilar morphologies [65] than with notable contributions from ion-beam fabricated copper-oxygen clusters due to the relative low concentration of copper in compare with that for carbon and other partial components (Table I). Also we have to note one more time that according to reported above core-level analysis (please see discussion above) not all the copper have interacted with the lattice oxygen, partially remaining in a metallic form, and as a distinct visually XPS spectral signature one can see the lower O 2s signal in the VB of C_x-TiO₂:[Cu⁺][Cu⁰] *thin-film* towards the *bulk* C_x-TiO₂: [Cu⁺][Cu²⁺]. At the same time a little forward, exactly this contingently "the lack of copper-oxygen interactions" *thin-film* sample exhibiting a little-

bit lower intensities and band-shape transformations in the O 2p region of XPS VB base-band area (see Fig. 5, the BE-range from 7 to 5 eV). Such a mentioned behavior well coincides with the offered interpretation and above reported core-levels analysis.

The BE-vicinity of 10 eV belongs to Ti 4s partial densities of states [44], and they are well-resolved by XPS in the untreated TiO₂-host employed as XPS external standard but not in TiO₂:[Cu⁺][Cu⁰] *thin-film* and *bulk* C_x-TiO₂:[Cu⁺][Cu²⁺] (Fig. 5). From Atomic Calculations of Photoionization Cross-Sections and Asymmetry Parameters [67] follows that the relation among Ti 4s, O 2p, and Cu 3d cross-sections σ for Al K α X-ray excitation is $0.5 \cdot 10^{-3}/0.24 \cdot 10^{-3}/1.2$, respectively, so it is not a surprise that the Ti 4s states are not resolved in XPS VB spectra of carbon-copper modulated samples. Taking into account Ref. [58] and $\sigma(\text{Ti } 3d) = 0.17 \cdot 10^{-3}$ with a usual location at 2.2 eV and the relatively weak contribution of Ti 4s even into VB of TiO₂, then the suppression of Ti states in the valence band becomes evident. The VB Base-band Width (BBW) is enlarged greatly from 4.3 eV for untreated TiO₂ up to 7 eV because of the appeared Cu 3d contribution. With that, the BBW becomes of Cu 3d majority character with all further consequences for the material electronic properties. Earlier the dominating character of Cu 3d electronic states in the BBW region was established independently employing X-ray emission overlay approach for the triple-component oxide [68], thus proving the suppressing character of copper 3d stated appearance in BBW. No essential influence of carbon on the electronic structure of final carbon-copper modulated samples were found except C 1s core-level transformations.

Table II. Formation energies (in eV per defect) for various combinations of defects – oxygen vacancies (vO), substitutional carbon impurities (C_S), substitutional (S) and interstitial (I) copper ions – in different TiO_2 matrices. The most probable configurations of defects in realistic matrices (see discussion in the text) are marked with bold font.

Host/Impurity	rutile (bulk)	anatase (bulk)	anatase (surface)
C	+8.77	+8.58	+3.18 in bulk -0.26 on surface
vO	+4.43	+4.57	+3.31
vO after C_S	+3.95	+3.89	+2.90
1S	+6.54	+6.06	+4.14
C_S	+5.10	+5.81	+6.13
vO	+4.41	+4.67	+0.27
$C_S + vO$	+4.28	+4.22	+0.15
2S	+6.40	+5.84	+4.29 (for 3S +4.45)
C_S	+6.15	+4.56	+3.90
vO	+4.48	+4.11	+1.92
$C_S + vO$	+4.20	+3.97	+1.85
S+I	+2.54	+2.33	+2.08
C_S	+2.10	+1.74	+1.98
vO	+1.51	+1.46	+0.41
$C_S + vO$	+1.39	+1.36	+0.97
2S+I	+3.27	+3.13	+2.82
C_S	+2.95	+2.84	+2.60
vO	+1.63	+1.55	+2.25
$C_S + vO$	+1.55	+1.49	+1.98

In order to understand and discuss the results of XPS measurements on a microscopic background, we performed the set of model calculations. First of all we examined the positions of carbon impurities and their role in the formation of oxygen vacancies. Results of our calculations, presented in Table II, demonstrate the large magnitude of the formation energy for the substitutional carbon impurity both in rutile and anatase *bulk* TiO_2 . In the vicinity of the surface, the energy required for incorporation of the carbon impurity is decreased but still remains valuable. In the case of placement of carbon atom on the surface as ad-atom (Fig. 6a) the formation energy turns to negative. Thus we can conclude that implanted carbon atoms should migrate towards the surface, being driven by the large difference in the formation energies for the *bulk* and for the vicinity of surface, with the onward yield out of TiO_2 grains and posterior formation of separate carbon phase. This result is in a

good agreement with performed XPS structure mapping which as well evidences about the formation of graphite-like phase (Fig.2).

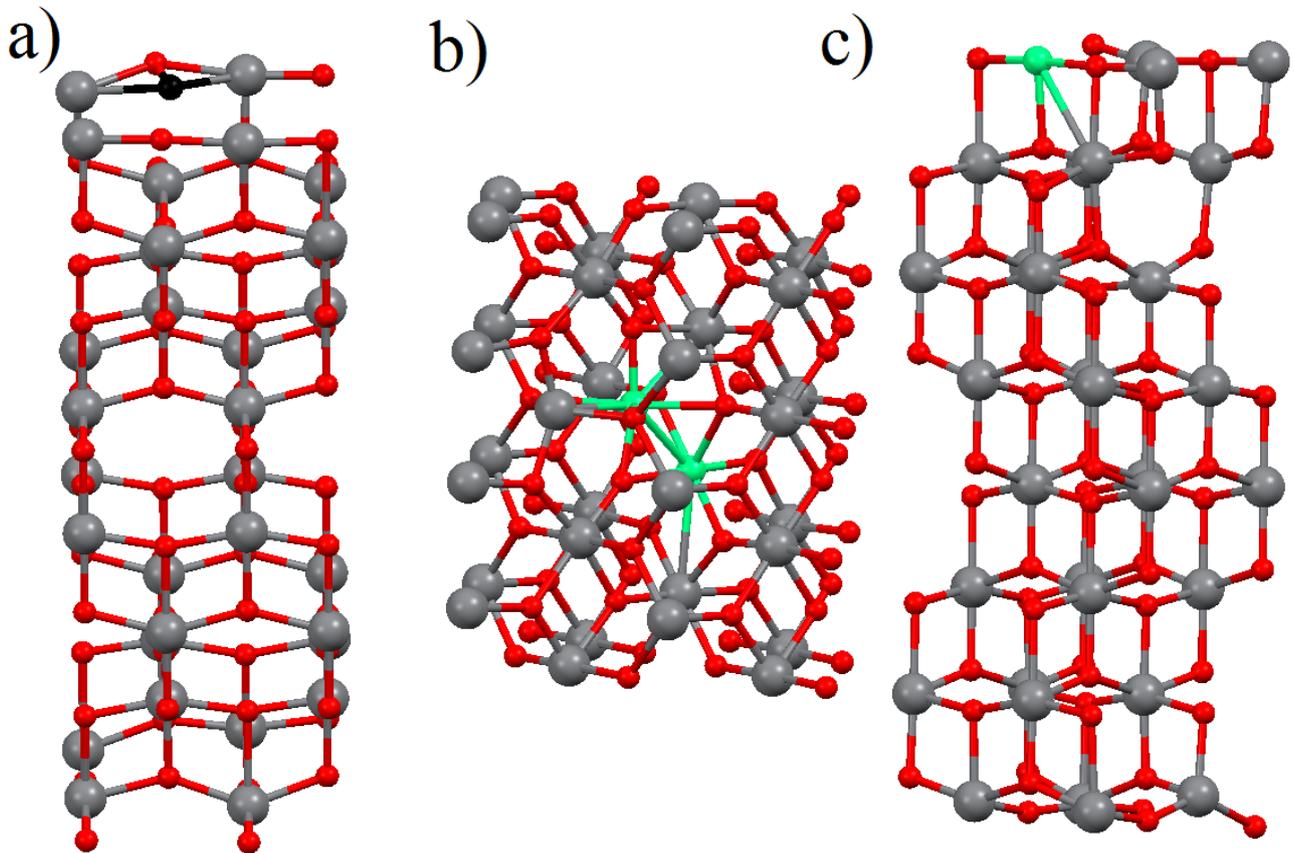


Figure 6. Optimized atomic structures of carbon ad-atoms: (a) on the surface (*thin-film*) of anatase TiO₂; (b) substitutional and interstitial copper impurities (S+I) in the *bulk* rutile TiO₂; (c) the combination of single substitutional copper impurity (1S) and oxygen vacancy on the surface of anatase TiO₂.

The formation energies of the oxygen vacancies both in rutile and anatase phases of TiO₂ are smaller than the energy cost of carbon implantation, but remain relatively high (see Table II). The energy difference between the defects in the *bulk* and the surface (*thin-film*) is much smaller than for the case of carbon impurities in the same positions. Calculations of the various configurations of the paired carbon impurity plus oxygen vacancy demonstrate that the lowest total energy of the system is

corresponding to the closest distance between them. Therefore carbon impurities can be discussed as the centers of the formations of all configurations for copper impurities in TiO_2 , i.e. new phase precursors. The performed calculation of the formation energies demonstrates that the presence of substitutional carbon impurity provides visible decreasing of formation energy of oxygen vacancy. Since the difference between formation energies of oxygen vacancies in the *bulk* and on the surface (*thin-film*) is not very large, so the migration of vacancies from bulk to the surface is much less favorable than for carbon impurities. This result is also in agreement with experimentally obtained XPS data, well evidencing the abundance of oxygen vacancies both in *bulk-rutile* and *thin films-anatase* after sequential carbon-copper loadings (Fig. 4).

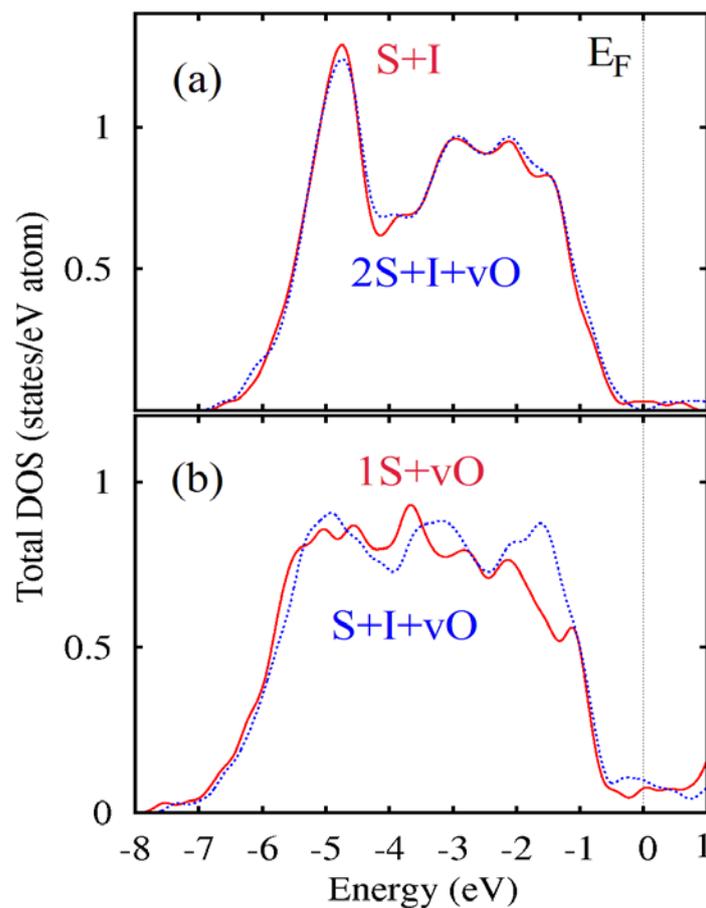


Figure 7. Base-band width Densities of States (DOSes) for the most energetically favorable configuration of copper defects and oxygen vacancies (a) in the *bulk* and (b) *thin-films* of TiO_2 .

The next step of our modeling is the analysis of energetics of various copper defects in TiO₂ matrices with and without defects. The results of calculations (see Table II) demonstrate the insignificant difference between rutile and anatase in the *bulk* morphology. For both phases of TiO₂ the preferable configurations of copper impurities in the absence of other defects are the combinations of substitutional and interstitial impurities (Fig. 6b). The nature of this phenomenon is the combination of larger ionic radii and lower valences of copper than titanium that makes simple substitution energetically much less favorable. The presence of other defects provides decreasing of the formation energies for all configurations of copper defects without varying the tendency to fabricate the clusters of substitutional and interstitial copper impurities. The Base-band area electronic structure shown in Fig. 7a demonstrates that the formation of discussed clusters provides clearly visible shift of the valence bands toward Fermi level. The presence of oxygen vacancies provides at the same time insignificant changes in electronic structure though. So the results of calculations are in agreement with XPS VB structure mapping which demonstrates the presence of metallic copper in the *bulk* rutile TiO₂ host-matrix.

In the case of a surface without defects, the results of our calculations demonstrate the same tendency for the cluster formation (see Table II) that are the combination of substitutional and interstitial defects, as it has been established in the *bulk* TiO₂. The presence of carbon impurities does not change valuable energetics of the copper defects configurations, in contrast with oxygen vacancies which create some hollow-space and make the formation of the single substitutional copper defect (Fig. 6a) also rather possible. The Base-band electronic structure (Fig. 7b) of the most probable defects demonstrates the smearing of the valence band in the *thin-film* C_x-TiO₂:Cu. The nature of this smearing is the presence of various types of oxygen atoms (surface, vicinity sub-surface, the environmental atoms of copper impurities, etc.) that provides the broadening of O 2p band (see Fig. 5 and 7a). At the same time the formation of copper clusters provides the regular increasing of the closest to Fermi level DOS'es. This result is also in qualitative agreement with experimental results

(Figs. 3 and 5) that demonstrate a combination of copper-oxide and metallic copper on the surface of oxygen-deficit TiO₂.

6. Conclusions

The heavy carbon-loadings with posterior Cu-sensitization of the *bulk* and *thin-film* TiO₂ hosts were made using sequential pulsed ion-implantation technique. The complete XPS qualification of the electronic structure of final C_x-TiO₂:Cu samples allows to establish the actual formulas as C_x-TiO₂:[Cu⁺][Cu²⁺] for the *bulk* and C_x-TiO₂:[Cu⁺][Cu⁰] for *thin-film* morphologies, respectively. No copper-carbon and titanium-carbon interactions were experimentally found due to the remaining majority of neutral C–C bonds (sp³-type) after carbon-copper implantation and only a lack of embedded carbon is fabricating the O–C=O clusters. This conclusion well agrees with XPS VB BBW analysis where it have been established the absence of BBW-broadening toward Fermi level. With that, there are no doubts about the dominating majority of Cu 3d states in comparance with other XPS BBW components, especially in the valence base-band of the *bulk* C_x-TiO₂:Cu. The performed theoretical model-scenarios for several trends of over the threshold carbon-loadings and posterior Cu-sensitization well agrees with experimentally mapped electronic structures of final C_x-TiO₂:Cu samples. It was shown that neutral carbon low-dimensional impurities play the role of precursors for the new phases while Cu-sensitizing the C_x-TiO₂ intermediate-state hosts.

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